

Three-body resonances from Lattice QCD

ω -meson, PRL 133 (2024) 21, arXiv:2407.16659, $\pi(1300)$: arXiv:2510.09476

Carsten Urbach

Gefördert durch



Deutsche
Forschungsgemeinschaft

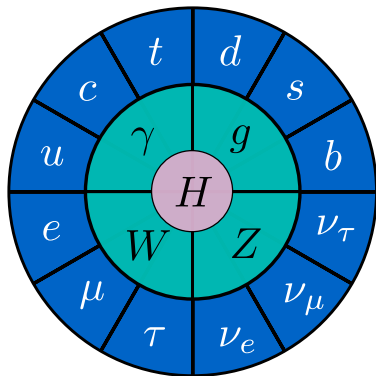
<NUMERIQS>



color meets flavor

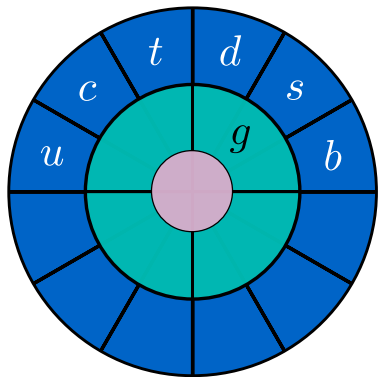


Standard Model of Particle Physics



- standard model highly successful

Standard Model of Particle Physics



- standard model highly successful
- strong interaction part
⇒ quantum chromodynamics (QCD)
- gives rise to **hadron spectrum**
mesons and baryons
- still new hadrons are being found
 X, Y, Z states
- **most hadrons are resonances**

The ω -meson

- experimentally discovered in 1961

[Maglich et al., PRL 7 (1961)]

- theoretically predicted by Nambu (1957) to explain the electromagnetic form factors of the nucleon

[Nambu, Phys. Rev. 106, 1366 (1957)]

- also expected in the vector meson dominance (VMD) picture

[Sakurai, Ann. Phys. 11, 1 (1960)]

- dominates the isoscalar response in VMD

1961 press conference on ω discovery



[NARA & DVIDS Public Domain Archive, photographer: D. Cooksey]

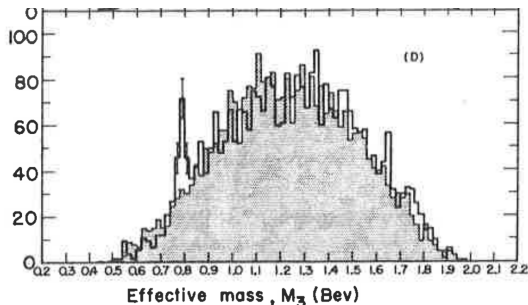
The ω -meson

- triplets identified in

$$\bar{p} + p \rightarrow \pi^+ + \pi^+ + \pi^- + \pi^- + \pi^0$$

decays

- $I^G(J^{PC}) = 0^-(1^{--})$
- PDG: $M_\omega = 782$ MeV, width 8.5 MeV
- $\approx 90\%$ decays to $\pi^+\pi^-\pi^0$
- about 1.5% $\rightarrow \pi^+\pi^-$



[B. Maglich et al, PRL 7 (1961)]

(BeV is one Bevatron \equiv GeV)

What Do We Know Theoretically?

- in 2022, we reviewed the theory status of hadron resonances

[Mai, Meißner, Urbach, Phys.Rept. 1001 (2023) 1-66]

- focus on calculations from first principles

5. Results: Well separated resonances.....
 - 5.1. The $\rho(770)$ -resonance
 - 5.2. The $K^*(892)$ -resonance
 - 5.3. The $\Delta(1232)$ -resonance
 - 5.4. The $f_0(500)$ -resonance
6. Results: Coupled channels/thresholds.....
 - 6.1. Light mesons
 - 6.2. The Roper-resonance $N(1440)$
 - 6.3. Specific open and closed charm systems.....
 - 6.3.1. The $D_0^*(2300)$
 - 6.3.2. The $D_0^*(2300)$ from experimental data .
 - 6.3.3. The $D_{s0}^*(2317)$
 - 6.3.4. The $X(3872)$
 - 6.3.5. The $Z_c(3900)$
 - 6.4. Other exotic states.....
 - 6.4.1. States involving heavy-light mesons
 - 6.4.2. Dibaryon states.....

What Do We Know Theoretically?

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- focus on calculations from first principles

given the huge number of experimentally known states:

surprisingly little is known theoretically!

almost exclusively $H \rightarrow XY$

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 - $\rho(3872)$
 - $\rho(3900)$
 - states.....
 - involving heavy-light mesons
 - on states.....

The $\pi(1300)$ Excited Pion

- first excited pion state

- first observed in 1981

[Ananeva et al., JETP Lett. 34, 488 (1981)]

- mass of 1300 ± 100 MeV
200 – 600 MeV width
- even existence debated

[P. d'Argent et al., JHEP 05, 143]

Citation: R.L. Workman et al. (Particle Data Group), Prog.Theor.Exp.Phys. **2022**, 083C01 (2022)

$\pi(1300)$ DECAY MODES

Mode		Fraction (Γ_i/Γ)
Γ_1	$\rho\pi$	seen
Γ_2	$\pi(\pi\pi)_{S\text{-wave}}$	seen
Γ_3	$\gamma\gamma$	

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- almost degenerate to the $\eta(1295)$ and $K(1460)$
- interesting insights into emergence of hadrons

⇒ even for excited pions room for significant contributions!

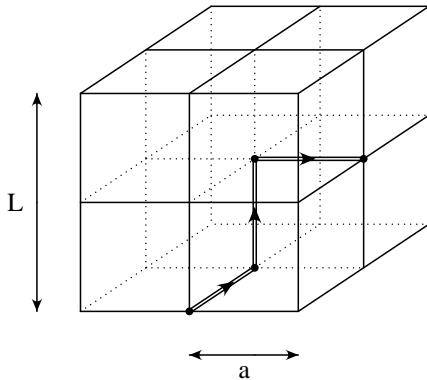
Quantum Chromodynamics

$$S[A_\mu, \bar{\psi}, \psi] = \int d^4x \left\{ \frac{1}{4} G_{\mu\nu}^2 + \bar{\psi}_q (i\gamma_\mu D_\mu + m_q) \psi_q \right\}$$

- astonishingly simple action, intriguingly complex dynamics
 - running coupling: QCD is non-perturbative at low energies
- ⇒ hadron spectrum requires non-perturbative methods

Lattice QCD Regularisation

- quantum field theory requires regularisation
- lattice regularisation:
 - ⇒ discretise space-time
 - hyper-cubic $L^3 \times T$ -lattice with lattice spacing a
 - ⇒ momentum cut-off: $k_{\max} \propto 1/a$
 - derivatives \Rightarrow finite differences
 - integrals \Rightarrow sums
 - gauge potentials A_μ in $G_{\mu\nu} \Rightarrow$ link matrices U_μ (' $\bullet \longrightarrow \bullet$ ')
- work in Euclidean space-time \Rightarrow **use Monte Carlo**



Lattice QCD Regularisation

- Monte Carlo: access to equilibrium, vacuum properties
- fundamental observables:
Euclidean correlation functions

$$\langle \mathcal{O}_i^\dagger(p, t) \mathcal{O}_j(p, t') \rangle \propto \sum_n c_{i,n} c_{j,n} e^{-E_n t}$$

- with interpolating operators \mathcal{O}_i with certain quantum numbers
- simulations at bare parameters need to renormalise

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- **continuum limit:**

$$\lim_{a \rightarrow 0}$$

(i.e. at least 3 lattice spacing values)

- **infinite volume limit:**

$$\lim_{L \rightarrow \infty}$$

- **physical mass limit:**

$$\lim_{m_\ell \rightarrow m_\ell^{\text{phys}}} \quad \text{or} \quad M_\pi^2 \rightarrow (M_\pi^{\text{phys}})^2$$

And then: Compute the Spectrum?

- lattice stochastic methods:
work in finite volume / Euclidean space-time
- ⇒ real valued, quantised eigenvalues of the lattice Hamiltonian
no continuum of states
- Maiani and Testa:
interactions properties cannot be studied directly
[\[Maiani and Testa, \(1990\)\]](#)
- ⇒ there is no one-to-one correspondence of an energy level to a resonance state
- the connection is only provided by the Lüscher method!

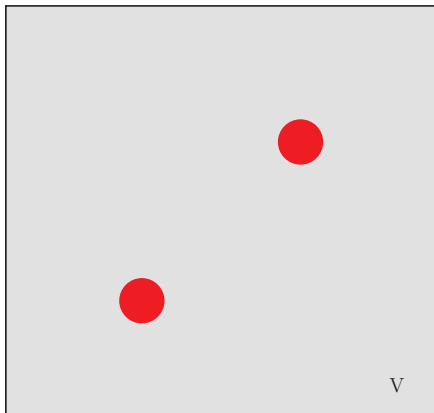
And then: Compute the Spectrum?

yes and no!

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Lüscher Method

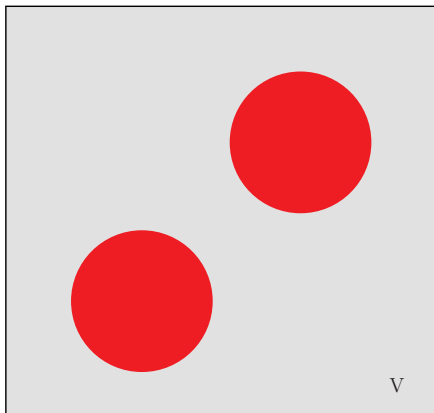
finite volume: boon and bane!



- for $V \rightarrow \infty$:
 - \Rightarrow interaction probability very low
 - $\Rightarrow E_{2p}(p=0) = 2E_{1p}$

Lüscher Method

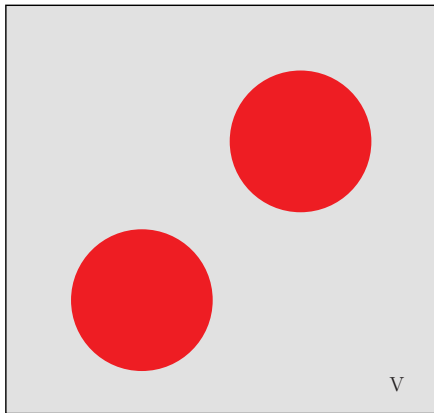
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- Lüscher: correction in $1/V$ related to scattering properties!

The 1 + 1-dimensional Analog

- plane wave acquires phase shift $\delta(k)$
- finite extend L , periodic BC

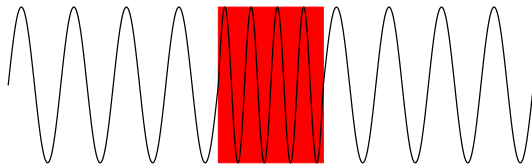
$$e^{ikL+2i\delta(k)} = e^{ik0} = 1$$

- quantisation condition

$$k_n L + 2\delta(k_n) = 2n\pi$$

- momenta k_n from dispersion relation

$$W_n = 2\sqrt{m^2 + k_n^2}$$

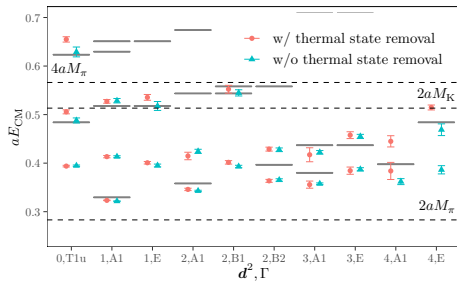


Procedure

- 1 determine non-interacting m
- 2 determine energies W_n
- 3 $W_n \rightarrow k_n$
- 4 $k_n \rightarrow \delta(k_n)$

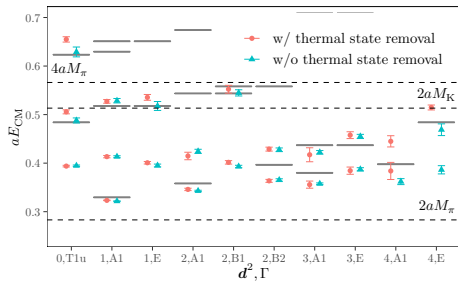
The General Case

Lattice Energy Levels E



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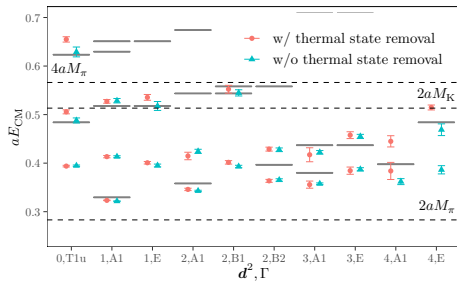
Determinant Equation

$$\det [\mathcal{M}^{\Gamma, \mathbf{d}}(E) - \cot(\delta)]^\Lambda = 0$$

(\mathcal{M} Lüscher function)

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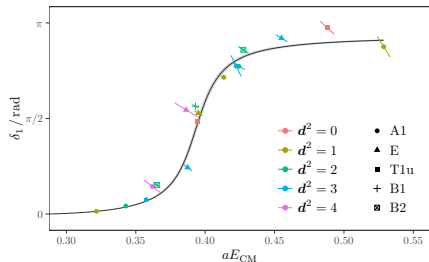
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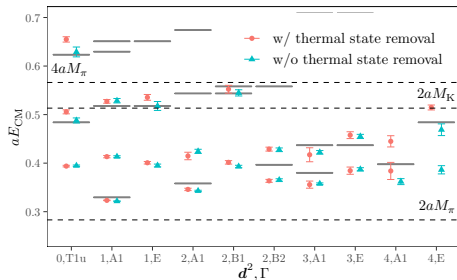
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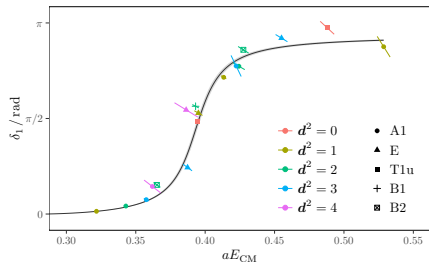
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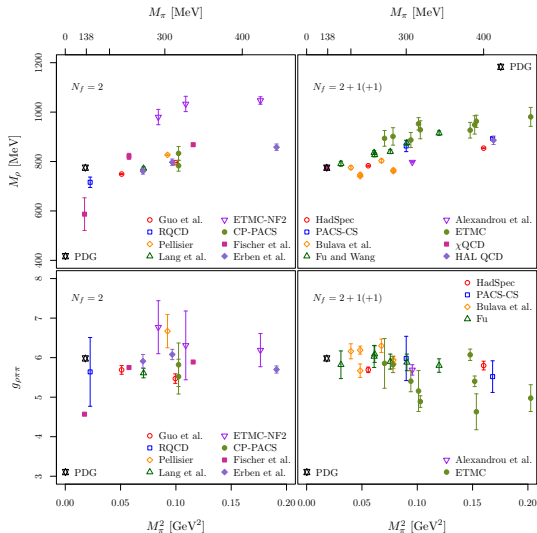


Parametrise Energy Dependence
e.g. Breit-Wigner or better



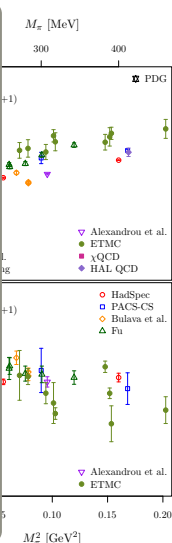
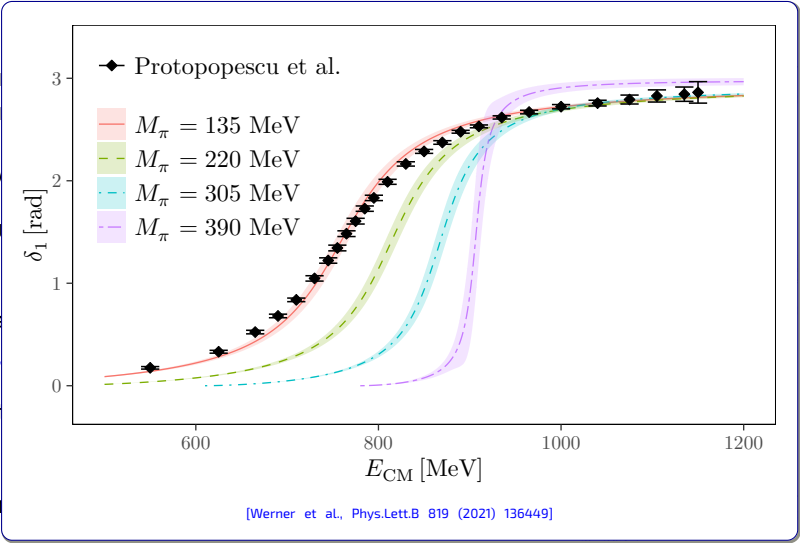
Example: The $\rho(770)$ Meson

- ρ -resonance a poster Breit-Wigner resonance
- ρ decays to $\pi\pi$ in a P-wave
- best studied resonance from Lattice QCD
- summary of 16 Lattice studies
[Mai, Meißner, Urbach, Phys. Rept. 1001 (2023) 1-66]
- bare lattice results for $N_f = 2$ and $N_f = 2 + 1(+1)$
- systematics clearly visible



Example: The $\rho(770)$ Meson

- ρ -reso
- resona
- ρ decay
- best st
- QCD
- summa
- [Mai, Meißner,
- bare la
- $N_f = 2$
- system



A Third Particle Enters the Game

- Three particle decays highly relevant
- Three-pion decays of
 $K, \eta, \omega, a_1(1260), a_1(1420), \pi(1300)$
- Decays of exotica, e.g.:
 $X(3872) \rightarrow \bar{D}^* D \rightarrow \bar{D} D \pi,$
 $Y(4260) \rightarrow J/\psi \pi \pi$
- Roper resonance $\rightarrow \pi N$ and $\pi \pi N$

Lattice Energy Levels E
Finite Volume, discrete, real

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Lattice Energy Levels E
Finite Volume, discrete, real



EFTs

Interaction Properties
Infinite Volume, possibly complex

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Three equivalent EFTs

- RFT (Relativistic Field Theory)
[Hansen, Sharpe, 2014]
- NREFT (Non-relativistic EFT)
[Hammer, Pang, Rusetsky, 2017]
- FVU (Finite Volume Unitarity)
[Mai, Döring, 2017]

Here: Finite Volume Unitarity (FVU)

builds on

- unitarity (\rightarrow probability conservation)
- analyticity (\rightarrow causality)
- crossing symmetry

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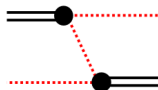
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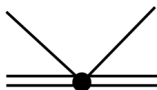
Quantisation Condition

$$\det \left[2L^3 E_{\mathbf{p}} (\tilde{K}^{-1} - \Sigma^L) - B - C \right]^{\Lambda} = 0$$

[Mai & Döring, EPJA 53 (2017)]



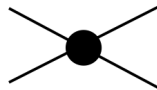
B



C



Σ



\tilde{K}^{-1}

Lattice Details

- based on CLQCD ensembles

[CLQCD, Hu et al, PRD 109 (2024)]

- tadpole tree-level Symanzik improved gauge
- $N_f = 2 + 1$ dynamical quark flavours
- tadpole improved clover fermions
- one level of stout smearing with $\rho = 0.125$
[Morningstar and Peardon, PRD 69 (2004)]
- possible $O(a\alpha_s)$ lattice artefacts

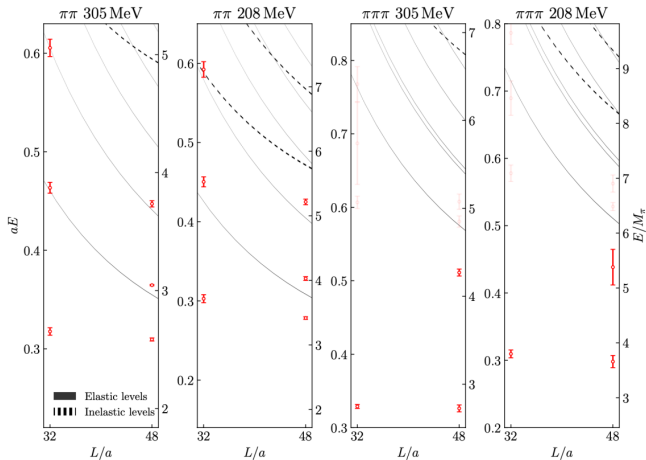
Ensemble	Volume	M_π/MeV	N_{confs}
F32P21	$32^3 \times 64$	206.8(2.1)	459
F48P21	$48^3 \times 96$	207.58(76)	221
F32P30	$32^3 \times 96$	303.61(71)	777
F48P30	$48^3 \times 96$	304.95(49)	201

- two pion mass values
- with two volumes each
- $a = 0.07746(18)$ fm

Finite Volume Spectrum: Parametrizations

First look at energy levels:

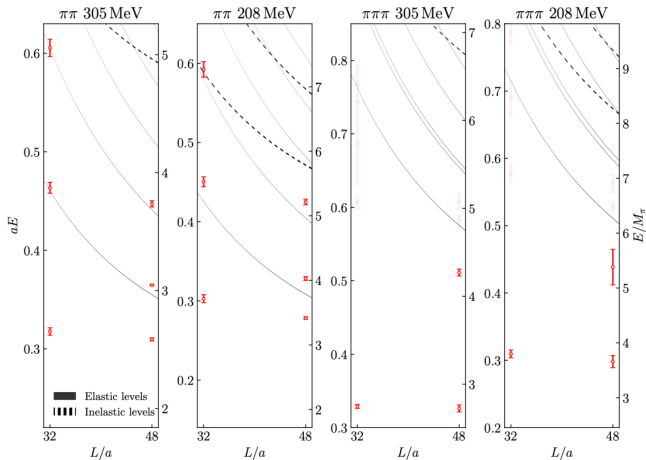
- red: interacting energy levels
- see attractive interaction for $\pi\pi$ and $\pi\pi\pi$
- $M_\pi = 305$ MeV: ω bound
- $M_\pi = 208$ MeV: ω decays
- need to parametrise



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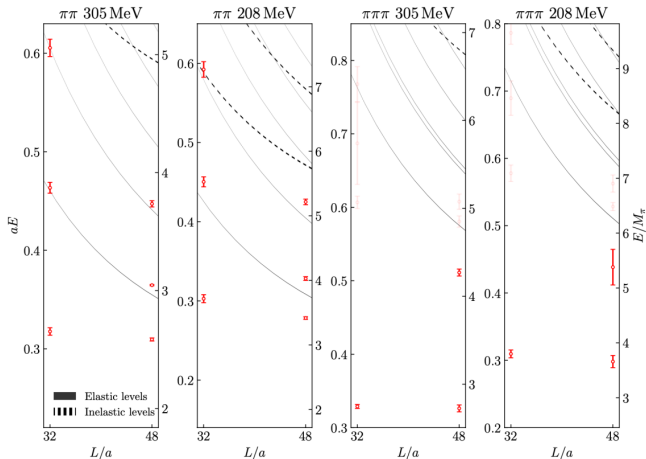
$$\det \left[2L^3 E_{\mathbf{p}} (\tilde{K}^{-1} - \Sigma^L) - B - C \right]^{\Lambda} = 0$$

Generic Form

$$K^{-1} \propto a_0 + a_1 \sigma_p(s)$$

$$c_{11} = \frac{c_0}{s + M_{\omega}^2} + c_1$$

σ_p : two-body invariant mass



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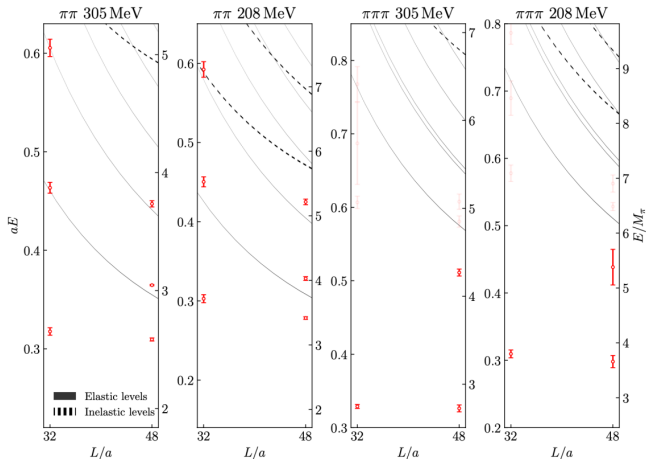
EFT Form

- EFT with vector mesons

[for a review: Meißner, Phys.Rept. 161 (1988)]

- re-express in $M_{\pi,\rho\omega}$, f_{π} and g

$$g_{\rho\pi\pi} = g_{\omega\rho\pi} = g$$



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- $f_\pi(M_\pi^2)$ from chiral PT

[Gasser,Leutwyler, Annals Phys. 158 (1984)]

EFT2 Form (two parameters)

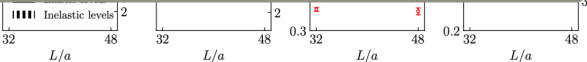
- $M_\rho = \sqrt{2} g f_\pi$ (KFSR)
- $M_\omega = M_\rho + \delta$

EFT4 Form (four parameters)

- $M_\rho = M_V + d M_\pi^2$
- $M_\omega = M_\rho + \delta$

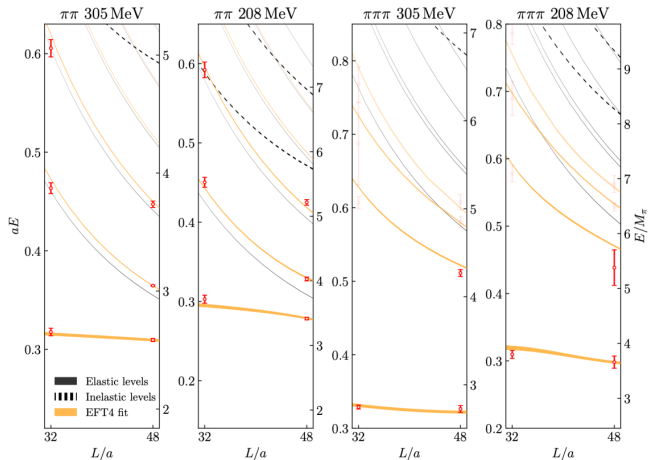
aE

E/M_π



Finite Volume Spectrum: EFT4 Fit

- parameters g, M_V, d, δ
- data described reasonably well
- even not included energy levels described
- removing ensemble with smallest $M_\pi \cdot L$ doesn't change results significantly
- $\chi^2/\text{dof} = 2.3$



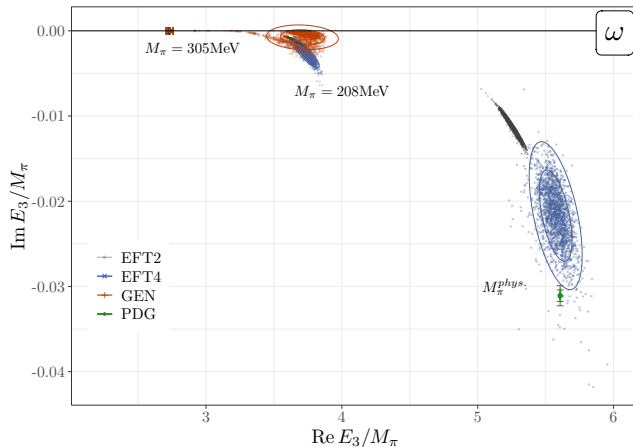
Final Step: Pole Positions (ω meson)

- with C, \tilde{K} as input
- solve an integral equation
[Mai, Döring, EPJA 53 (2017)]
- analytically continue to 2nd Riemann Sheet

⇒ Pole positions of ω

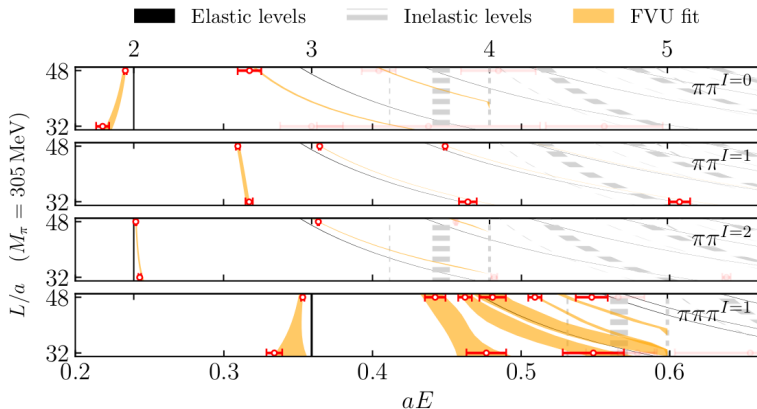
$$778.0(11.2) - i3.0(5) \text{ MeV}$$

[Yan et al. PRD, PRL 133 (2024) 21]



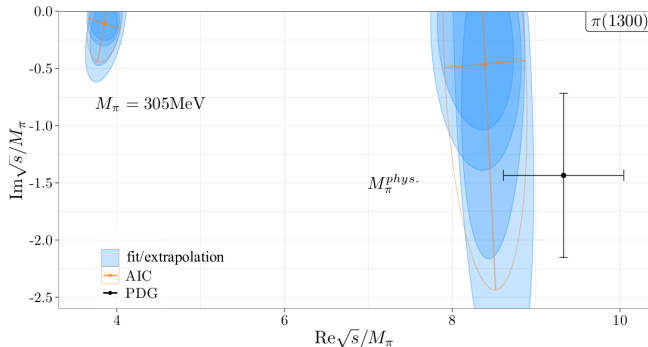
$\pi(1300)$: The Same Procedure

- more difficult due to large expected resonance mass
- three $\pi\pi$ isospin channels relevant
- ground state is a single pion
- for $M_\pi = 305$ MeV sufficient no of states below $6M_\pi$
- at $M_\pi = 208$ MeV too few states



$\pi(1300)$ Pole Position

- rely on $M_\pi = 305$ MeV ensembles only
- perform a model averaging procedure
- inverse amplitude method for extrapolation in M_π^2
- with large uncertainties we find evidence for a pole



The (notorious) Challenge: Extracting Energy Levels

- Euclidean Correlation Functions

$$C(t) = \sum_{l=0}^{N_s} A_l e^{-E_l t}$$

$$E_0 < E_1 < \dots$$

- obtained from stochastic (MC) simulations
- want to estimate energy levels E_l
- and amplitudes A_l (matrix elements)

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Challenge

- signal-to-noise problem

$$\text{StN} \propto \exp(-\Delta E t)$$

with in general $\Delta E > 0$

[Lepage (1989)]

- signal deteriorates exponentially
- increasing severity with l

Example: The Effective Mass of the Nucleon

- effective mass

$$M_{\text{eff}}(t) = -\frac{1}{\delta t} \log \left(\frac{C(t + \delta t)}{C(t)} \right)$$

- since $E_0 < E_{l \neq 0}$

$$\lim_{t \rightarrow \infty} M_{\text{eff}}(t) = E_0$$

- all other contributions to C
exponentially suppressed
- but t finite
 \Rightarrow excited state contaminations

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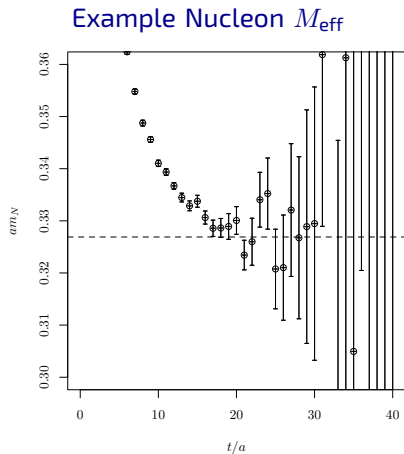
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[ETMC (2021)]

Example: The Effective Mass of the Nucleon

- effective mass

$$M_{\text{eff}}(t) = -\frac{1}{\delta t} \log \left(\frac{C(t + \delta t)}{C(t)} \right)$$

- since $E_0 < E_{l \neq 0}$

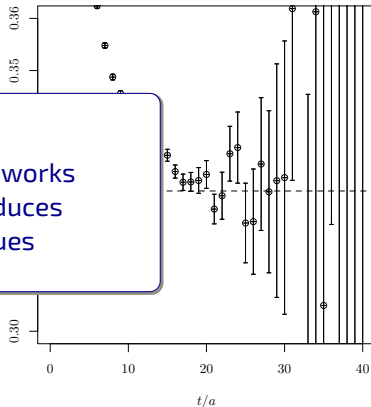
$$\lim_{t \rightarrow \infty} M_{\text{eff}}(t) = E_0$$

- all other contributions to C exponentially suppressed
- but t finite
 \Rightarrow excited state contaminations

Goal

Find a method that either works at smaller t -values or reduces the noise at large t -values

Example Nucleon M_{eff}



[ETMC (2021)]

Generalised Effective Mass

- build correlator matrix

$$C_{\alpha\beta}(t) = \langle O_{\alpha}^{\dagger}(0) O_{\beta}(t) \rangle$$

- O_{α} with appropriate quantum numbers
- cast as eigenvalue problem

$$C(t_0)^{-1} \cdot C(t) v_n(t, t_0) = \lambda_n(t, t_0) v_n(t, t_0)$$

- with gen. effective mass

$$\lambda_n(t, t_0) \propto e^{-E_n(t-t_0)}$$

- but: no. of operators limited...

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- but: no. of operators limited...

- create more operators for free!

$$O_{\Delta t, \alpha}(t) = O_{\alpha}(t + \Delta t) = T(\Delta t) O_{\alpha}(t)$$

Generalised Pencil of Function method

[Schiel, PRD92, 034512 (2015)]

- leads to the square Hankel matrix

$$H_{ij}^{\alpha\beta, m}(t) = C_{\alpha\beta}(t + i + j), \quad i + j < 2m$$

- and the eigenvalue problem

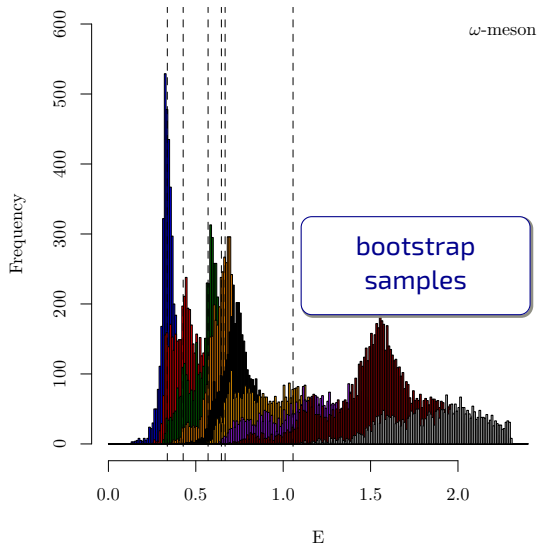
$$H^m(t_0)^{-1} \cdot H^m(t) v_n(t, t_0) = \lambda_n(t, t_0) v_n(t, t_0)$$

And Here Comes the Noise...

- the method enforces an effective noise model

$$e^{\text{Re}(E)t} e^{i \text{Im}(E)t}$$

- the larger the matrix, the more spurious modes are found
- and they mix with the physical ones



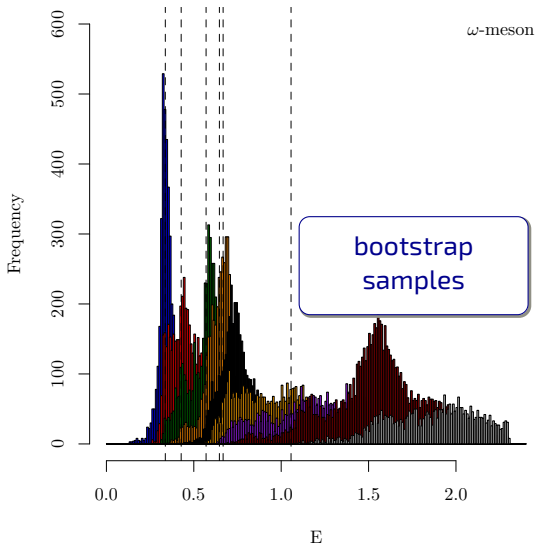
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find/invent criteria for physical states
or reformulate the problem!



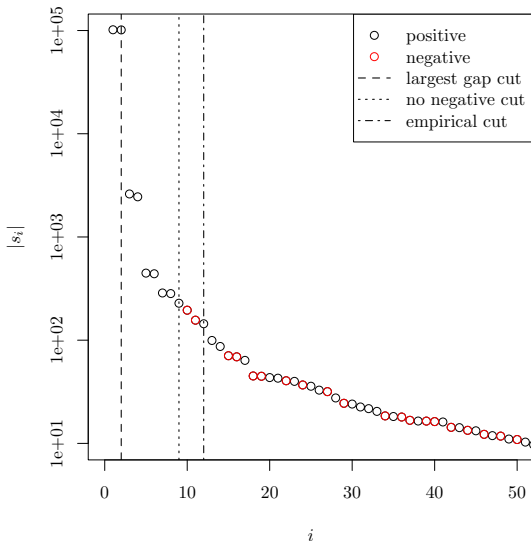
The Reformulated Problem

What is the best approximation to the correlator with k terms in a Frobenius norm?

- solution is known!
- diagonalise H (or SVD)

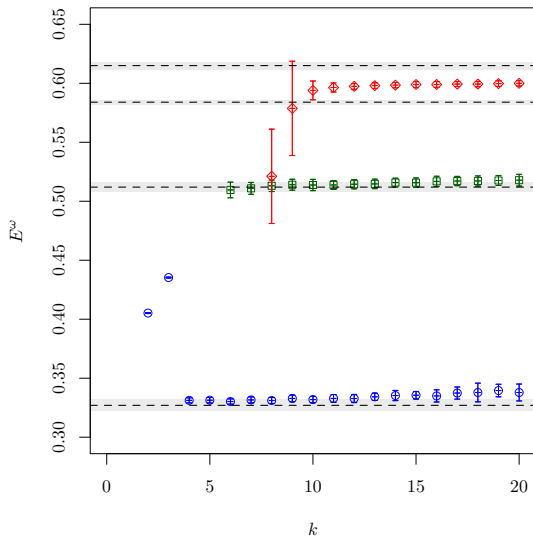
$$H = U^\dagger S U$$

- set all but the k largest modulus eigenvalues $|s_{i>k}|$ to zero
- reconstruct $\tilde{H} = U_k^\dagger S_k U_k$ and solve for eigenvalues



THC Results

- truncated Hankel correlator (THC) method
[\[Ostmeyer, Urbach, arXiv:2510.15500\]](#)
- single parameter k
- all data included
- stable spectrum for $k \geq k_{\text{cut}}$
- directly applicable to **imaginary time evolution** in quantum computing



What can be improved...

- ω and $\pi(1300)$ are clearly exploratory studies!
- more ensembles, more M_π -values, larger volumes
- more energy levels by including more frames and irreps
- include also operators with strange quarks
- investigate the φ and the ω - φ mixing
- take the continuum limit
- include other partial waves

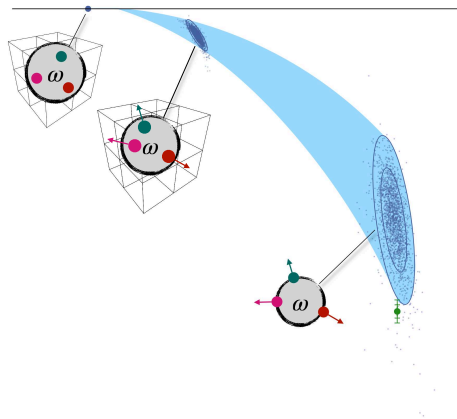
Summary

- resonances in LQCD challenging problem
- presented a first calculation of the ω and $\pi(1300)$ resonances
- using EFTs, obtained results at physical point
- ω pole position

$$(778(11) - i3.0(5)) \text{ MeV}$$

- $\pi(1300)$ pole position

$$(1169(46) - i62(170)) \text{ MeV}$$



[Yan, Mai, Garofalo, Meißner, Liu, Liu, Urbach, arXiv:2407.16659]

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Marco Garofalo



Johann Ostmeyer