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The Dawn of Nuclear Supremacy: First Bounds on Ultralight Dark Matter from the Thorium Isomer Transition

Eric Madge (IFT-UAM/CSIC)

Implications of the laser excitation of the Th-229 nucleus for dark matter searches

Fuchs, Kirk, **EM**, Paranjape, Peik, Perez, Ratzinger, Tiedau; [arXiv:2407.15924](https://arxiv.org/abs/2407.15924) [hep-ph]

Introduction

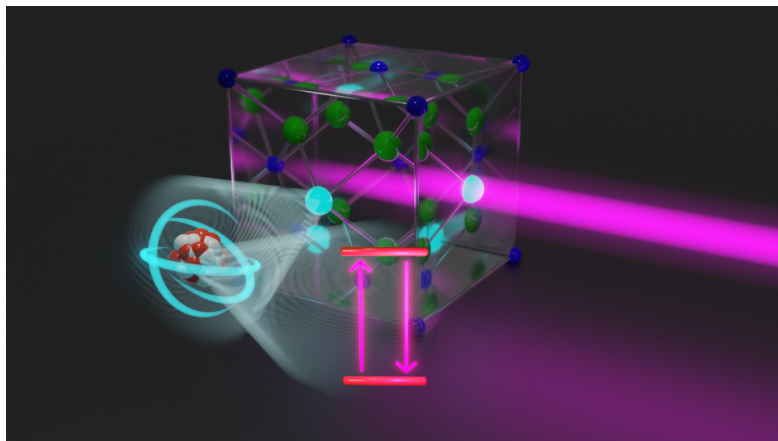
first laser excitation of
 $^{229\text{m}}\text{Th}$ in crystals:

■ CaF_2 : Tiedau et al., PRL 132 (2024)

$$f = 2020.409(7) \text{ THz}$$

■ LiSrAlF_6 : Elwell et al., PRL 133 (2024)

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credit: TU Wien/PTB

⇒ breakthrough on the path towards $^{229\text{m}}\text{Th} \rightarrow ^{229}\text{Th}$ nuclear clock

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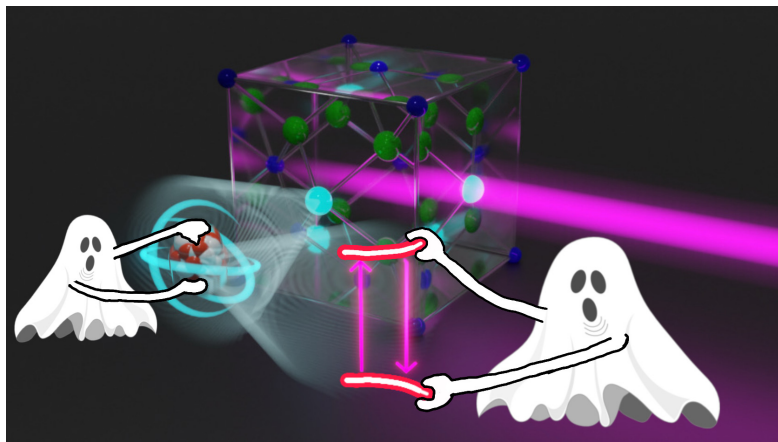
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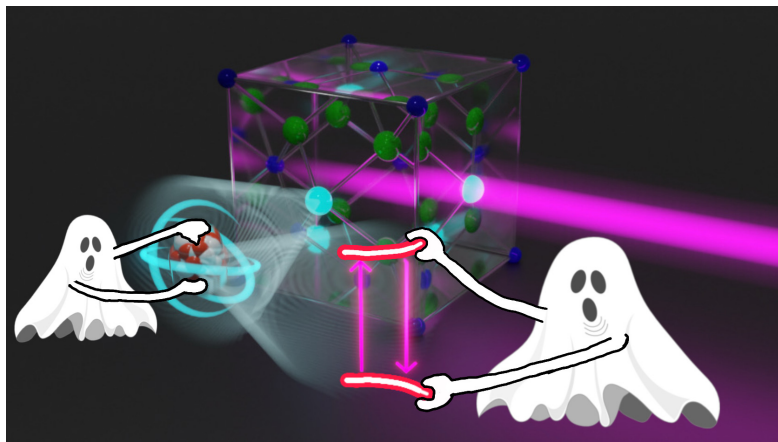
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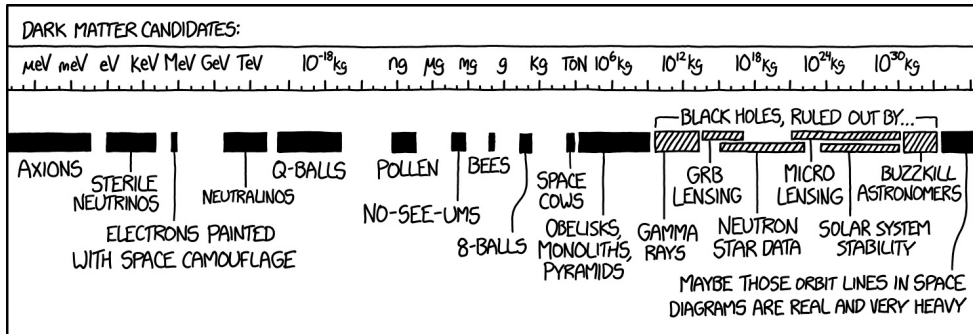


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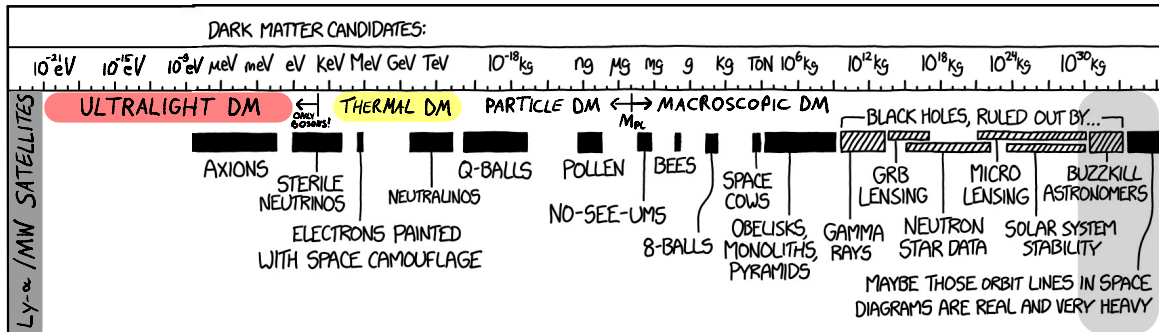
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Ultralight Dark Matter



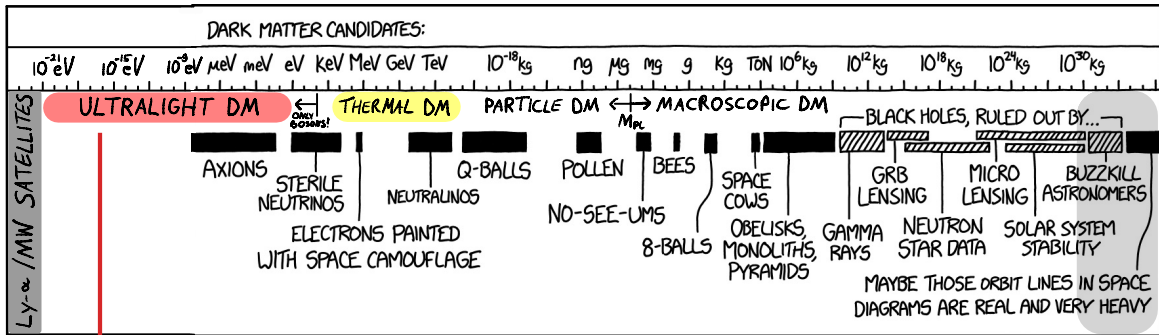
credit: <https://xkcd.com/>

Ultralight Dark Matter



credit: <https://xkcd.com/> [edited]

Ultralight Dark Matter



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- behaves like classical field

$$\langle N \rangle \sim n \lambda_{\text{dB}}^3 \sim \frac{\rho}{m} (mv)^{-3} \gg 1 \quad \Rightarrow \quad m_{\text{DM}} \lesssim 1 \text{ eV}$$

- oscillates coherently ($\tau_{\text{coh}} \sim \frac{2\pi}{mv^2} \sim 10^6$ oscillations)

$$\ddot{\phi} + 3H\dot{\phi} + m^2\phi^2 = 0 \quad \Rightarrow \quad \phi(t) \sim a^{-\frac{3}{2}} \cos(mt + \delta)$$

Nuclear-Coupled Dark Matter

■ scalar ϕ :
$$\mathcal{L} \supset -\frac{\beta_s}{2g_s} \underbrace{\sqrt{4\pi G_N}}_{\kappa} d_g \phi G_{\mu\nu}^a G^{a\mu\nu}$$

$$\Rightarrow \Lambda_{\text{QCD}} = \Lambda_{\text{QCD}}^{(0)} (1 + \kappa d_g \phi)$$

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$$\Rightarrow m_\pi^2 \sim m_\pi^{2(0)} \left(1 - \frac{m_u m_d}{4(m_u + m_d)^2} \frac{a^2}{f_a^2} \right)$$

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$\Rightarrow$ oscillations in ^{229}Th frequency:

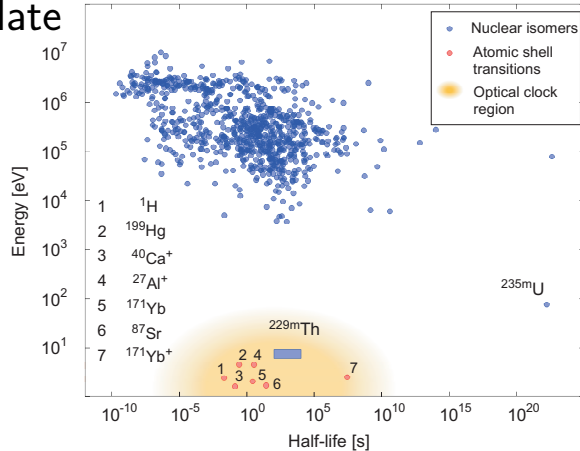
$$\frac{\Delta f_{\text{Th}}}{f_{\text{Th}}} \propto d_g \phi \propto \cos(m_\phi t) \quad \text{or} \quad \frac{\Delta f_{\text{Th}}}{f_{\text{Th}}} \propto \frac{a^2}{f_a^2} \propto \cos(2m_a t)$$

^{229}Th as a Nuclear Clock Candidate

$^{229\text{m}}\text{Th} \rightarrow ^{229}\text{Th}$ isomer transition:

■ lifetime: $\tau \approx 30$ min

■ energy: $\omega = 8.36$ eV (2020 THz)



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$$\omega = \Delta E_{\text{nuc}} + \Delta E_C$$

~ 8 eV $\Rightarrow$ cancellation
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nuclear energy

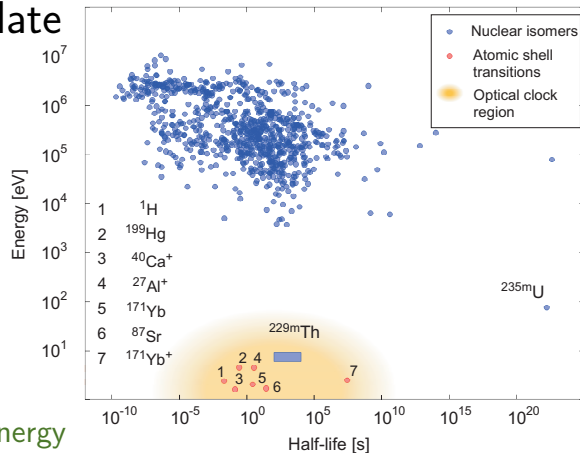
$\rightarrow \propto \Lambda_{\text{QCD}}$

$\rightarrow \sim \mathcal{O}(\text{MeV})$

Coulomb energy

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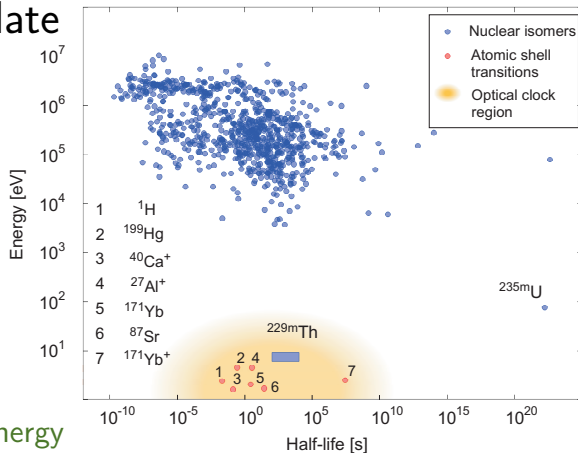
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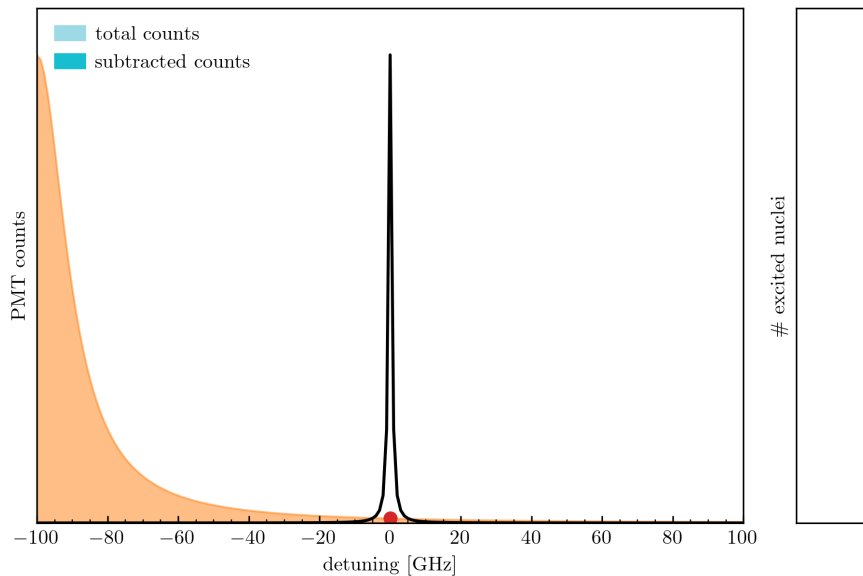
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$$\Rightarrow \text{enhanced sensitivity: } \frac{\Delta f}{f} \sim \frac{\Delta E_{\text{nuc}}}{\omega} \frac{\partial \log \Delta E_{\text{nuc}}}{\partial \log \Lambda_{\text{QCD}}} \frac{\Delta \Lambda_{\text{QCD}}}{\Lambda_{\text{QCD}}} \sim 10^5 \frac{\Delta \Lambda_{\text{QCD}}}{\Lambda_{\text{QCD}}}$$

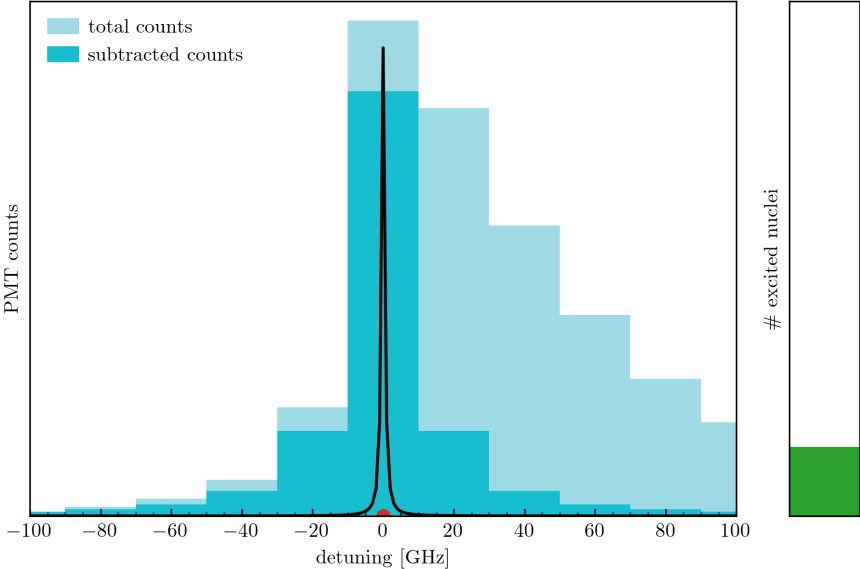


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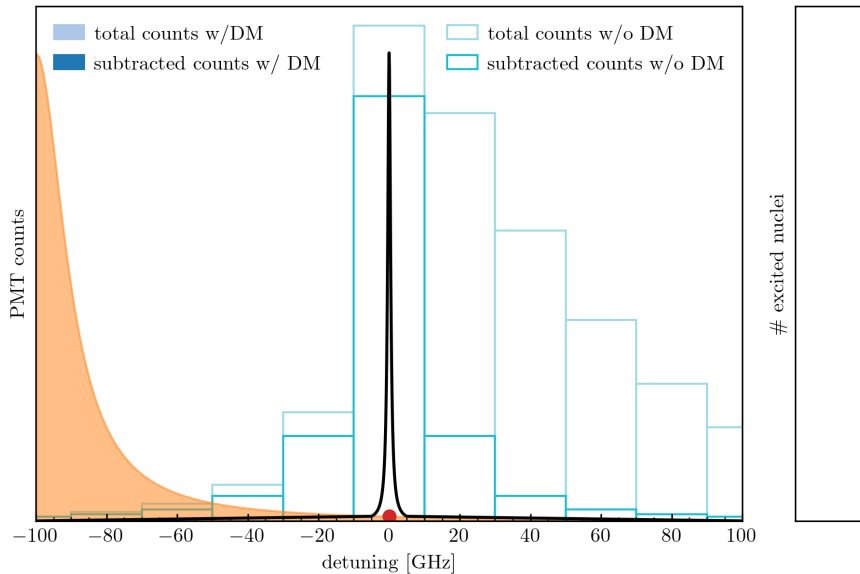
Excitation Spectrum



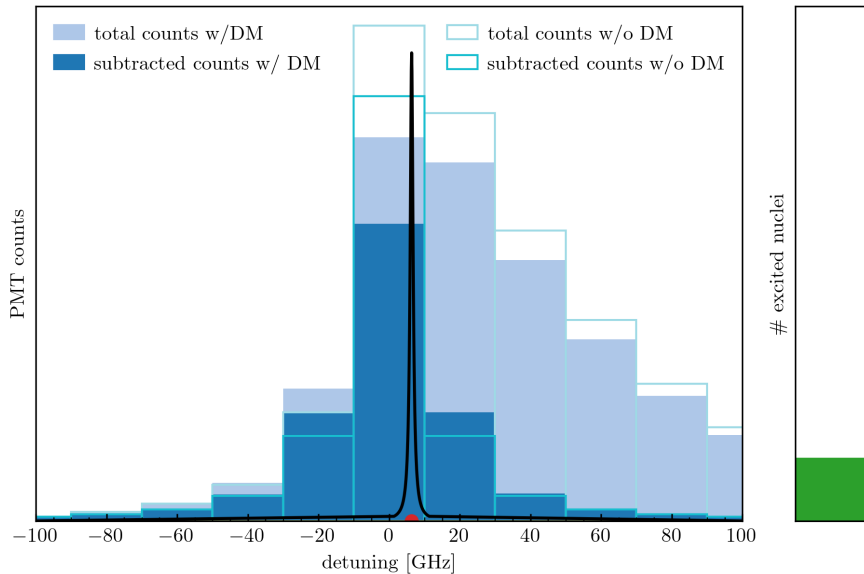
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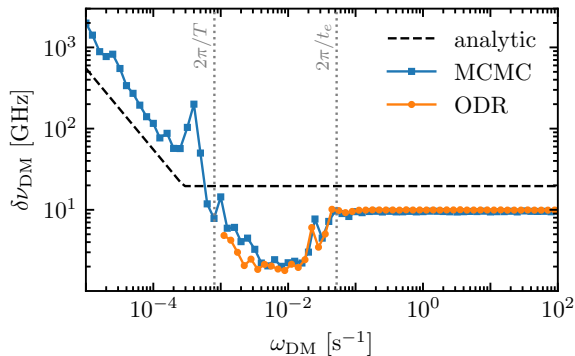


Lineshape Fit to Tiedau et al.

fit to excitation data at fixed ω_{DM}
to constrain δf_{DM}

■ $\omega_{\text{DM}}T \ll 1$: linear drift between scans

$$\delta f_{\text{DM}} \lesssim \frac{\sigma_{\text{Th}}}{\omega_{\text{DM}}T} \quad \sigma_{\text{Th}} - \text{uncertainty}$$



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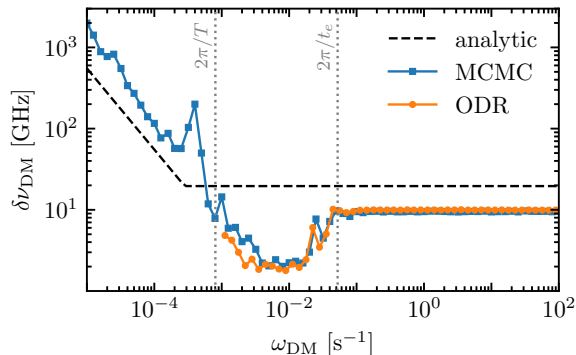
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- $\omega_{\text{DM}}(t_e + t_d) \gg 1$: broadening by $2\delta f_{\text{DM}}$

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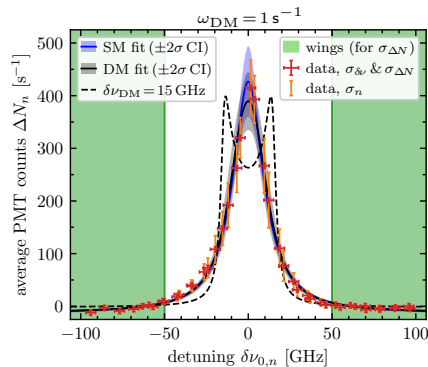
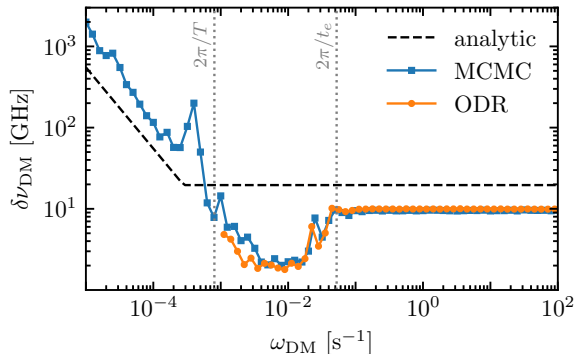
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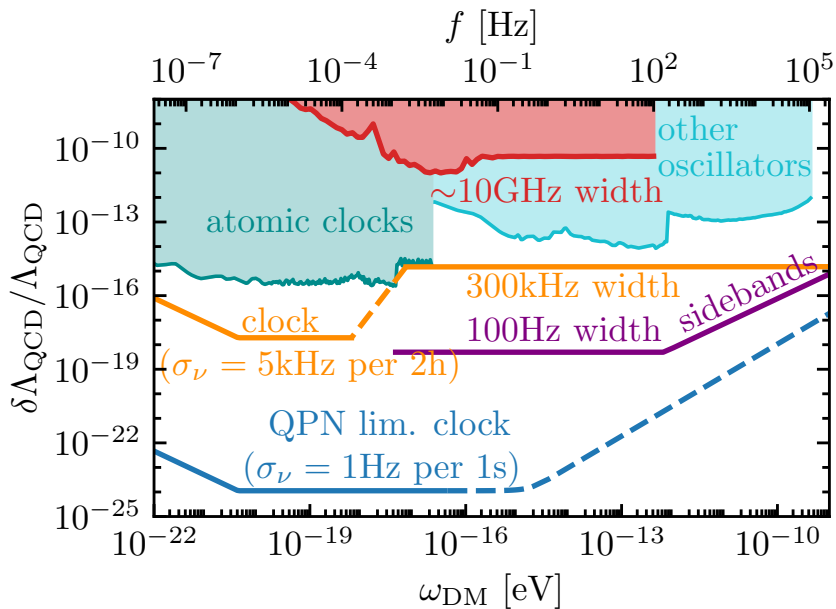
$$\delta f_{\text{DM}} \lesssim \Delta f \quad \Delta f - \text{FWHM}$$

■ $\delta f_{\text{DM}} \gg \omega_{\text{DM}}$: two peaks at $f_{\text{Th}} \pm \delta f_{\text{DM}}$

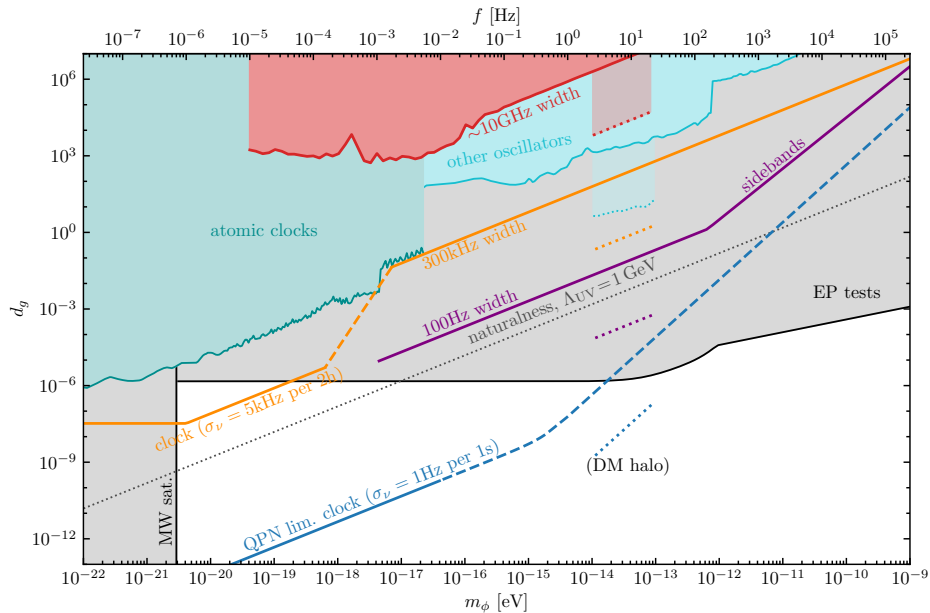
$$\langle I(f) \rangle \propto \frac{\theta \left(1 - \left| \frac{f - f_{\text{Th}}}{\delta f_{\text{DM}}} \right| \right)}{\sqrt{\delta f_{\text{DM}}^2 - (f - f_{\text{Th}})^2}}$$



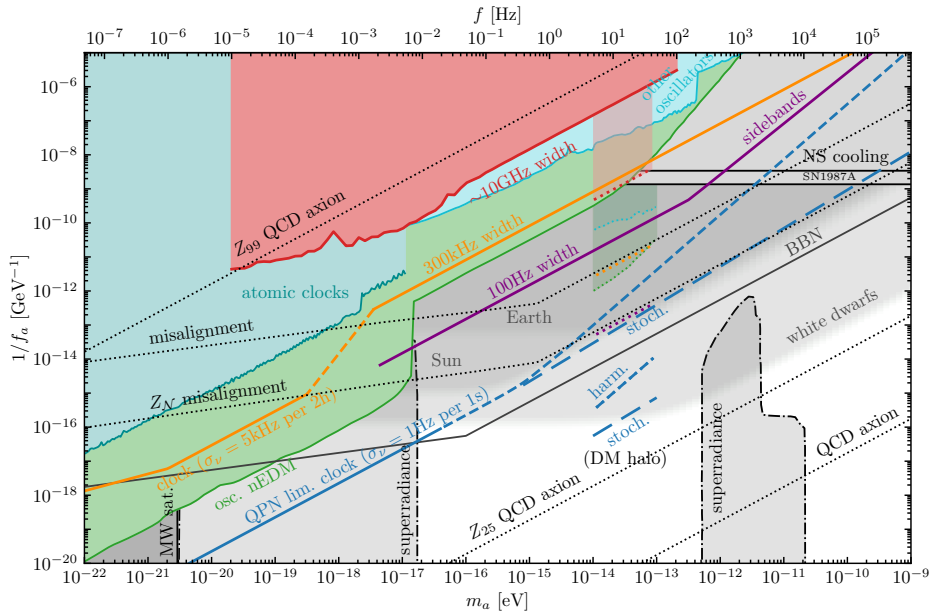
Bounds on Variations of Λ_{QCD}



Scalar Dark Matter Bounds



Axion Dark Matter Bounds



Conclusion

- ULDM can induce oscillations in fundamental constants
- quantum clocks can be used to probe ULDM
- nuclear transitions can be used to probe ULDM even before proper clock operation
- current technology can already probe new parameter space

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Backup

Excitation Lineshape Model

detection efficiency $\rightarrow$ $\Delta N_n = \varepsilon \int_0^{t_e} dt' \Gamma(t') \int_{t_e}^{t_e+t_d} dt \frac{1}{\tau} e^{-\frac{t-t'}{\tau}},$

excitation probability between t' and $t' + dt'$ $\rightarrow$ $\Gamma(t')$
 probability for nucleus excited at t' to decay between t and $t + dt$ $\rightarrow$ $\frac{1}{\tau} e^{-\frac{t-t'}{\tau}}$

laser lineshape $\rightarrow$ $L(f; f_L(t'))$ (Lorentzian)
 thorium lineshape $\rightarrow$ $I(f; f_{Th}(t'))$ ($\delta(f - f_{Th}(t'))$)

$\Gamma(t') \propto \int df L(f; f_L(t')) I(f; f_{Th}(t'))$

$$= \Delta N_{\text{off}} + \mathcal{N} \int_0^1 dx \frac{e^{-xt_e/\tau}}{1 + 4 \left(\frac{\delta f_n(x)}{\Delta_L} \right)^2}$$

count offset $\rightarrow$ ΔN_{off}
 prefactor $\rightarrow$ $\mathcal{N}$
 laser linewidth $\rightarrow$ Δ_L

frequency detuning: $\delta f_n(x) = \delta f_{L,n} + \delta f_{\text{off}} + \delta f_{\text{DM}} \cos(\omega_{\text{DM}} t_e x - \varphi_n)$

DM modulation $\rightarrow$ δf_{DM}
 $f_{\text{Th}} - f_{L,n}$ $\rightarrow$ $\delta f_{L,n}$
 detuning offset $\rightarrow$ δf_{off}