Holographic timelike entanglement entropy and the black hole interior

Based on [hep-th: 2408.15752] and work in progress, with M. P. Heller and A. Serantes

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New insights in black hole physics from holography Madrid, 17 June 2025



Introduction and motivation

Entanglement entropy

From quantum mechanics to holography

• $|\Psi\rangle$ pure state. Select a region A, then the reduced density matrix $\rho_A = \text{Tr}_{\bar{A}} |\Psi\rangle\langle\Psi|$. Its von Neumann entropy is the entanglement entropy: Sorkin'83

$$\mathbf{S}_{\mathbf{A}} \coloneqq -\mathrm{Tr}_{\mathbf{A}}(\rho_{\mathbf{A}}\log\rho_{\mathbf{A}})$$

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• Given a (d - 1)-dimensional region A, find a codimension-2 extremal surface γ_A in the bulk spacetime, anchored on ∂A .

$$S_A \coloneqq \min_{\operatorname{Area}[\gamma_A]} \frac{\operatorname{Area}[\gamma_A]}{4G_N^{(d+1)}}$$

Ryu, Takayanagi '06; Hubeny, Rangamani, Takayanagi '07 **1**/7

Bulk reconstruction and new holographic probes

• Emergence of spacetime: entanglement wedge reconstruction.

Czech, Karczmarek, Nogueira, Van Raamsdonk '12; Headrick, Hubeny, Lawrence, Rangamani '14; Dong, Harlow, Wall '16



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 Emergence of spacetime: entanglement wedge reconstruction.



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• Especially rare are observables that probe the **black hole interior**.



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 What can we say on one-sided probes of the black hole interior? The holographic dual to timelike entanglement

How do we define timelike entanglement?

Analytic continuation of the spacelike result

• CFT_2 result and its analytic continuation as $\Delta x^2 - \Delta t^2 < 0$:

Doi, Harper, Mollabashi, Takayanagi, Taki '23





NB: in this talk, we will be agnostic on the general field theory definition of timelike entanglement.

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Geometric interpretation of timelike entropy

Complex bulk extremal surfaces

- Consider a (d 1)-dim. timelike region A.
- Find a complex codimension-2 extremal surface γ_A in a complexified bulk, anchored on ∂A (real boundary). Then:





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- Hatches with analytic continuation of all known closed form expressions.
- + Can be applied to **any spacetime**.



Timelike entanglement and the black hole interior

Timelike entanglement in SAdS₄

• Strip subsystem *A*. Metric:

$$ds^{2} = \frac{1}{z^{2}} \left(-f(z) dt^{2} + \frac{dx_{\parallel}^{2}}{f(z)} + d\mathbf{x}_{\perp}^{2} \right)$$
$$f(z) = 1 - \left(\frac{z}{z_{H}}\right)^{d}$$



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• New: multiple surfaces!

 Vacuumdisconnected ones probe a complexified black hole interior.



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Interior from analytic continuation

Heller, FO, Serantes '25 (to appear)







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 New phase transition in spacelike entanglement! Slightly (δ) before crossing the lightcone:



Interior from analytic continuation

Heller, FO, Serantes '25 (to appear)



- **Dominant contributions** continue to timelike entropy.
- Subdominant contributions continue to the « BH interior ».

Slightly (δ) before crossing the lightcone:



Summary and outlook

- **Timelike entanglement entropy** is understood geometrically only in a **complexified bulk spacetime**.
- **One-sided** holographic probes that arise from a **phase transition in spacelike entanglement** can reach a **complexified black hole interior**.



Summary and outlook

- **Timelike entanglement entropy** is understood geometrically only in a **complexified bulk spacetime**.
- **One-sided** holographic probes that arise from a **phase transition in spacelike entanglement** can reach a **complexified black hole interior**.
- Relation to **two-sided settings**?
- Complex holographic probes: an (almost) unexplored landscape.
- *Field theoretical realisations?* Tensor networks, Gaussian fields.



Thank you!

Simplest check of the prescription: AdS₃

- Consider $A = \left[-\frac{1}{2}\Delta t, \frac{1}{2}\Delta t\right]$, at constant x = 0.
- We look for a geodesic joining two timelike-separated points at the boundary $z = \varepsilon \rightarrow 0$:

$$z(\lambda) = i \sqrt{\frac{1}{4}\Delta t^2 - \varepsilon^2} \operatorname{csch} \lambda, \qquad t(\lambda) = \sqrt{\frac{1}{4}\Delta t^2 - \varepsilon^2} \tanh \lambda$$

• The boundary conditions are at:

$$\lambda_m = \log \frac{\Delta t}{\varepsilon} + \frac{i\pi}{2} + O(\varepsilon)$$

Reλ

 $\boldsymbol{\chi}$

Z

A

 Δt

Simplest check of the prescription: AdS₃

• An interesting path:

along paths 1A and 1B: $z, t \in \mathbb{R}$; along path 2: $z, t \in i \mathbb{R}$.



t

 $\boldsymbol{\chi}$

Ζ

A

 Δt

Simplest check of the prescription: AdS₃

• Another example: along paths 1A and 1B: $z, t \in \mathbb{R}$; along paths 2A and 2B: z, t = const; along path 3: $z, t \in \mathbb{R}$.







Vaidya spacetime in 2+1 dimensions

• Metric and mass-function:

$$ds^{2} = \frac{-(1 - m(v) z^{2}) dv^{2} - 2dv dz + dx^{2}}{z^{2}}$$

$$m(v) = \frac{\alpha}{2}(1 + \tanh \gamma v)$$

• It represents the formation of a black brane of temperature $T = \sqrt{\alpha}/2\pi$ on a timescale $1/\gamma$ by the gravitational collapse of a shell of null dust.



An exactly solvable timelike quench

- **Thin-shell** limit in AdS₃-Vaidya: analytically solvable.
- The shell is defined by v = 0 (matching at real v). For a region $A = [t_C - \Delta t/2, t_C + \Delta t/2]$ we have:

$$S_{A}^{(T)} = i\pi + \begin{cases} 2\log(\Delta t), & t_{C} < -\Delta t/2 \\ 2\log\left[\frac{2}{r_{H}}\sinh\left(\frac{r_{H}(2t_{C} + \Delta t)}{4}\right) - \cosh\left(\frac{r_{H}(2t_{C} + \Delta t)}{4}\right)\left(t_{C} - \frac{\Delta t}{2}\right)\right] & -\Delta t/2 < t_{C} < \Delta t/2 \\ 2\log\left[\frac{2}{r_{H}}\sinh\left(\frac{r_{H}}{\Delta t}\right)\right] & t_{C} > \Delta t/2 \end{cases}$$

Balasubramanian, Bernamonti, Craps, Keränen,

Keski-Vakkurif, Müller, Thorlacius, Vanhoofd '12



 (\mathbf{T})

 Δt

 Δt

Crosscheck of complex geodesics



Timelike entanglement in SAdS₄

 Pick the surface with minimal real part of the area?

Prescription P1

$$S_{A}^{(T)} = \min_{\substack{\text{Re Area}[\gamma_{A}^{\mathbb{C}}]}} \frac{\text{Area}[\gamma_{A}^{\mathbb{C}}]}{4G_{N}^{(d+1)}}$$



t

 x_i

A

Timelike entanglement in SAdS₄

 $\mathcal{A}_{\mathrm{reg}} = \lim_{\epsilon o 0} \left(\mathcal{A} - rac{2}{\epsilon}
ight)$

- Problem: the UV/IR correspondence is violated!
- Possible resolution:

Prescription P2

$$S_A^{(T)} = \frac{\operatorname{Area}[\gamma_A^{\mathbb{C}}]}{4G_N^{(d+1)}}$$

for a $\gamma_A^{\mathbb{C}}$ such that **upon analytic** continuation $S_A^{(T)}$ reduces to standard entanglement entropy.



Towards a field theory definition

• Connection with « temporal entanglement » in tensor networks:



Bañuls, Hastings, Verstraete '09; Hastings, Mahajan '23; Carignano, Marimón, Tagliacozzo '23

Towards a field theory definition

• Tensor network evaluation of timelike entanglement entropy:



Carignano, Marimón, Tagliacozzo '23

 Can we do more? Quenches and time-dependent backgrounds can be explored with tensor networks.

Foligno, Zhou, Bertini '23