# **Krylov spread complexity** as holographic complexity beyond JT gravity

based on 2412.17785 with M.P. Heller and J. Papalini





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Bdy setup: Double-scaled SYK model

$$H_{\text{SYK}} = i^{p/2} \sum_{1 \le i_1 < \dots < i_p \le N} J_{i_1 \cdots i_p} \psi_{i_1} \cdots \psi_{i_p} \psi_{i_p} \cdots \psi_{i_p} \psi_{$$

quantum system of chord states  $|n\rangle$  governed by transfer matrix  $\hat{T}$ 

(N fermions with p-body interaction)  $V_{i_p}$ ,  $N, p \to \infty$  while  $\left|\log q\right| = p^2/N$  fixed

 $\Rightarrow$  Coupling averaged multi-point functions of  $H_{SYK}$  computable via auxiliary [Berkooz et al. '18]

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Bulk dual: Sine-dilaton gravity [Blommaert, Mertens, Yao, Papalini, Levine, Parmentier '23 '24 '25]

Ad: Thomas Mertens' talk on Thursday  $\int \mathscr{D}g \mathscr{D}\Phi \exp\left(\frac{1}{2 \log q}\right) d$ 

 $\Rightarrow$   $H_{grav}$  from canonical quantization of  $E_{ADM}$  coincides with DSSYK  $\hat{T}$ 

 $\Rightarrow$  Holographic dictionary relates bulk length with chord nr.  $\hat{L} = 2 \log q \hat{n}$ 

(N fermions with p-body interaction)  $V_{i_n}$ ,  $N, p \to \infty$  while  $\left|\log q\right| = p^2/N$  fixed

$$d^2x\sqrt{g}\left(\Phi R + \sin(2\Phi)\right) + bdy.$$

















Krylov spread complexity: Expectation value of the position of a state  $|\psi(t)\rangle$ 

$$C_{K}(t) = \sum_{n=0}^{\infty} n |\langle K_{n} | \psi(t) \rangle|^{2}, |\psi(t)\rangle$$

spreading over 'minimal basis of evolution / Krylov basis'  $|K_n\rangle$  [Balasubramanian et al. '22]

# $= e^{-iHt} |R\rangle, |K_n\rangle = orthonorm(\{H^n |R\rangle\})$



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KEY RESULT: Sine-dilaton length  $\langle \hat{L} \rangle$  equals DSSYK Krylov complexity  $C_K(t)_{\beta}$ 

# $\left[-\partial_{\Delta}(2pt.fn)\right]_{\Delta=0} = \langle \hat{L} \rangle = C_{K}(t)_{\beta}$

valid @finite temperature ( $\beta$ ) & (disk) quantum level (q) | [lliesiu, Mezei, Sarosi '22]

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KEY RESULT: Sine-dilaton length  $\langle \hat{L} \rangle$  $\langle \hat{L} \rangle = \left[ -\partial_{\Delta} (2pt \cdot fn) \right]_{\Delta=0} = -\partial_{\Delta} \left[ Z_{\beta}^{-1} - \partial_{\Delta} \left[ 2q^{-1} \right]_{\beta} \right]$  $= \left| \log q^{2} \right| \sum n \left| \langle n | Z \right|_{\beta} \right|$ @ $\beta \ge 0 \& 0 \le q \le 1 \mid (2\text{-par. gen. of [F]})$ 

# $= e^{-iHt} |R\rangle, |K_n\rangle = orthonorm(\{H^n |R\rangle\})$

equals DSSYK Krylov complexity 
$$C_K(t)$$
  
 $\frac{1}{L} = 0 |e^{-\tau \hat{H}_{grav}} e^{-\Delta \hat{L}} e^{-(\beta - \tau) \hat{H}_{grav}} |L = 0\rangle ]$   
 $Z_{\beta}^{-1/2} e^{-i\hat{T}(t - i\beta/2)} |0\rangle |^2 = |\log q^2 |C_K(t)_{\beta}$   
Rabinovici et al. '23] =  $(\beta \to 0, q \to 1)$  case





### Choice of reference state $|R\rangle$

Two inequivalent options for combined Lorentzian + Euclidean evolution

$$C_K$$
 for  $e^{-iT(t-i\beta/2)}|0\rangle$ 

 $\equiv |R\rangle$ 

$$\neq C_K \text{ for } e^{-iTt} e^{-\beta T/2} | 0 \rangle$$
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Two inequivalent options for combined Lorentzian + Euclidean evolution

$$C_{K} \text{ for } e^{-iT(t-i\beta/2)} | 0 \rangle \neq C_{K} \text{ for } e^{-iTt} e^{-\beta T/2} | 0 \rangle$$
$$\equiv | R \rangle \qquad \equiv | R \rangle$$

Starting from  $\langle L \rangle$  on the gravity side, we match this

# $\Rightarrow$ Gravity demands to assign complexity also to euclidean state preparation



Our result is valid for any  $q \Rightarrow$  Consistently, agrees with first quantum correction on top of the classical sine-dilaton gravity length [Bossi et al. '25] around  $q \approx 1$ 







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This talk:  $e^{-iT(t-i\beta/2)}|0\rangle$ 

### **Q: Precise influence of reference state choice on bulk manifestation?**

# $C_K$ specified by $H = T_{DSSYK}$ , $|R\rangle = |0\rangle$



This talk: 
$$e^{-iT(t-i\beta/2)}|0\rangle$$
  $C_K$  spe

### Q: Precise influence of reference state choice on bulk manifestation?

• Operator insertions in reference state:  $|R\rangle = O_{\Lambda} |0\rangle$  $\Rightarrow$  Dual manifestation as length in sine-dilation shockwave geometry?



# ecified by $H = T_{DSSYK}$ , $|R\rangle = |0\rangle$