

Entanglement and complexity in various braneworld models

Complexity: Between Field Theory and Gravity

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Based on work with Rob Myers, Shan-Ming Ruan and work in progress with Sergio Aguilar-Gutierrez, Ben Craps, Mikhail Khramtzov, Maria Knysh and Ashish Shukla

IFT Madrid, May 24 2023

Outline

- 1 Entanglement entropy and complexity
- 2 Entanglement and complexity of islands
 - Islands
 - Easy islands model
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- 3 Bounds from entanglement and complexity
 - Entanglement velocity and Lloyd bounds
 - Black holes + ETW branes
 - Bounds on intrinsic gravity on the brane
- 4 Conclusion

Geometry from entanglement

Entanglement entropy

Connection between entanglement and geometry

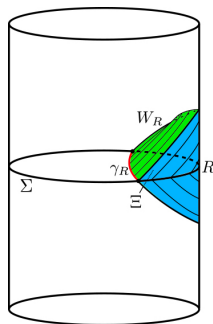
- HRT formula

$$S(R) = \min_{\partial\gamma_R = \partial R} \frac{A(\gamma_R)}{4G}$$

- EE from 1st law

$$\delta S = \delta\langle H \rangle \rightarrow E_{ab}^g = 0$$

- Entanglement wedge reconstruction

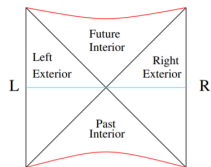


Geometry from entanglement

Entanglement entropy is not enough [Susskind]

Bulk regions inaccessible to entanglement entropy

- Entanglement entropy thermalizes fast
- HRT surfaces don't reach deep interior regions
- Deep interior regions not probed by entanglement entropy



Geometry from entanglement

Entanglement entropy is not enough

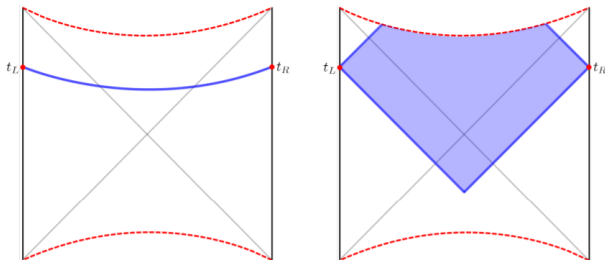
Solution: Study other geometric quantities which reach deeper into the bulk

Complexity

Holographic complexity

Two proposals: complexity=volume (C_V) and complexity=action (C_A)

$$C_V(R) = \max_{\partial\Sigma_R=R} \frac{V(\Sigma_R)}{G\ell}, \quad C_A(R) = \frac{I_{WdW}(R)}{\pi\hbar}$$

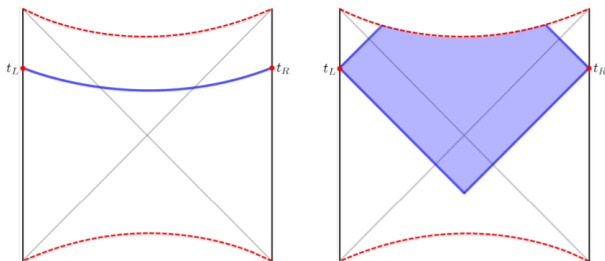


Complexity

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More proposals: CV2.0 and complexity = (almost) anything

Complexity

Circuit complexity

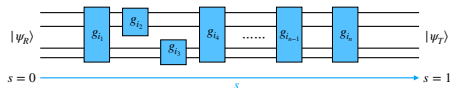
What is holographic complexity on the field theory side?

- Linear increase in time
- Slow to thermalize
- Reach very large values
- Switchback effect

State/circuit complexity

$$C_{\Psi_R}(\Psi_T) = \min_U D(U)$$

$$s.t. \quad U|\Psi_R\rangle = |\Psi_T\rangle$$

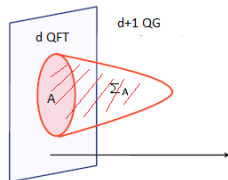


Complexity

Subregion complexity

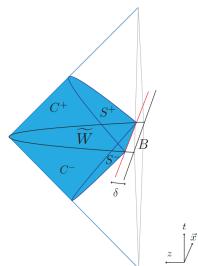
Considering subregion duality

- Holographic dual to reduced density matrix is its entanglement wedge
- Can we define the complexity of a subregion state?



Holographic subregion complexity

- Restriction of holographic proposals to entanglement wedge

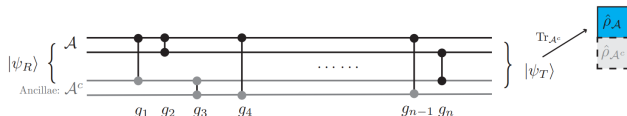
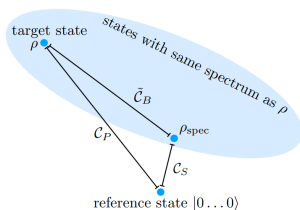


Complexity

Mixed state complexity [Agón, Headrick, Swingle]

What is holographic subregion complexity on the field theory side?

- Purification complexity
- Spectrum complexity
- Basis complexity
- ...



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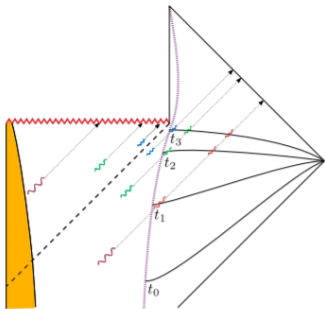
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Information paradox

Entropy curve

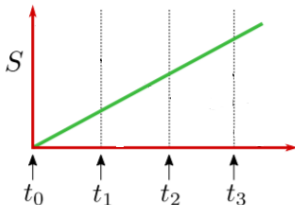
Hawking radiation

- Particle pair creation
- Near horizon
- Fall/escape of partners



Entanglement of black hole

- Pairs are entangled
- Constant Hawking radiation
- Linear increase in entropy

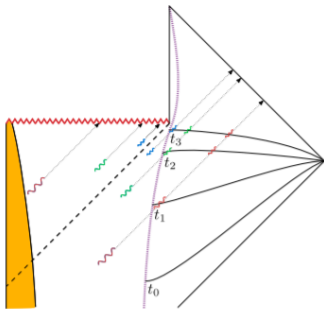


Information paradox

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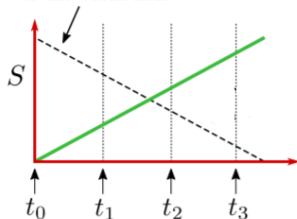
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Thermodynamic entropy
of the black hole



Generalized entropy and island rule

For a theory with gravity, generalized entropy

$$S_{gen}(R) = S_{EE}(R) + \frac{A(\partial R)}{4G}$$

In addition, the island rule

$$S_{gen}(R) = \min_I \left(S_{EE}(R \cup I) + \frac{A(\partial R)}{4G} + \frac{A(\partial I)}{4G} \right)$$

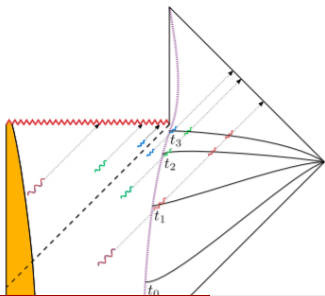
$\frac{A(\partial I)}{4G}$ is very big. Islands only relevant when there is a lot of entanglement between R and I

The island rule

[Almheiri, Mahajan, Maldacena, Zhao]

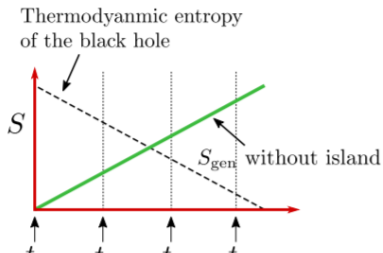
Island rule

- Generalized entropy
- When gravity is included
- Allow for “entanglement islands”



Page curve

- Early: no islands
- Entanglement increases
- Late: island configuration
- Entropy decreases

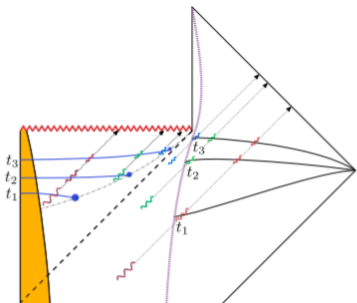


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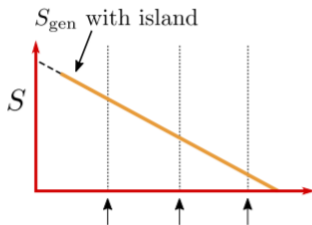
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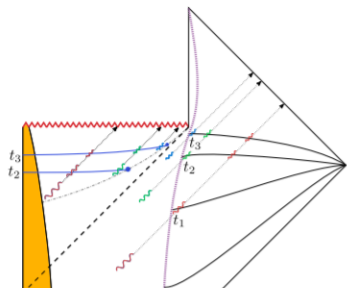


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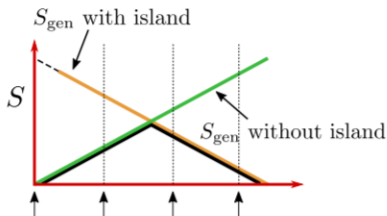
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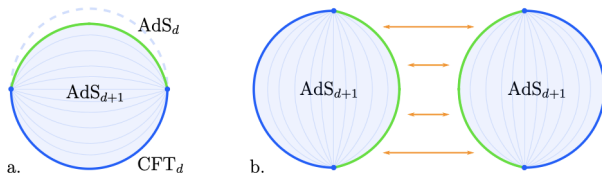
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QEs made easy

[Chen, Myers, Neuenfeld, Reyes, Sandor]

We use a Randall-Sundrum braneworld model



With action

$$I_{\text{bulk}} + I_{\text{brane}} = \frac{1}{16\pi G_N} \int d^{d+1}x \sqrt{-g} \left[R + \frac{d(d-1)}{L^2} \right] - T_0 \int d^d x \sqrt{-h}.$$

The location of the brane is determined by the Israel junction conditions

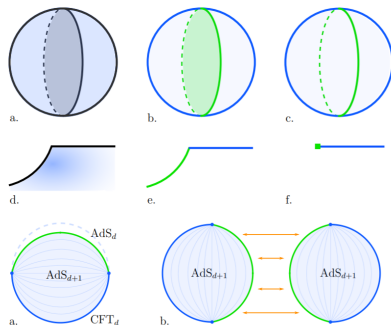
$$\Delta K_{ij} - h_{ij} \Delta K^k_k = -8\pi G_N T_0 h_{ij}.$$

QEs made easy

[Chen, Myers, Neuenfeld, Reyes, Sandor]

Two layers of holography, so three perspectives

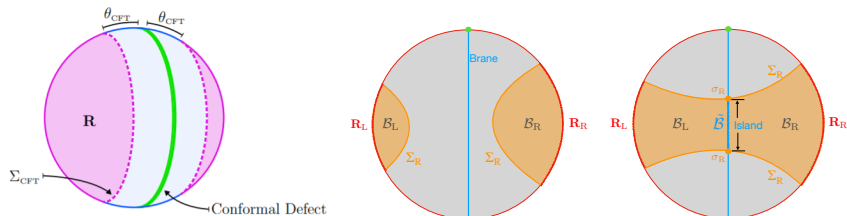
- Bulk: Asymptotically AdS_{d+1} + co-dimension 1 brane
- Brane: CFT in S_d coupled to brane with RS gravity + CFT
- Boundary: CFT in S_d with conformal defect



QEs made easy

Island phase transition

The island phase transition is due to the conventional transition found in holographic EE



Fefferman-Graham expansion

Ambient metric construction

$$ds^2 = g_{\mu\nu} dy^\mu dy^\nu = \frac{L^2}{z^2} (dz^2 + g_{ij}(z, x^i) dx^i dx^j)$$

Einstein equations fix g_{ij} in terms of boundary $g_{ij}^{(0)}$ (and $g_{ij}^{(d/2)}$)

$$g_{ij}(z, x^i) = g_{ij}^{(0)}(x^i) + \frac{z^2}{L^2} g_{ij}^{(1)}(x^i) + \dots + \frac{z^d}{L^d} \left(g_{ij}^{(d/2)}(x^i) + f_{ij}(x^i) \log\left(\frac{z}{L}\right) \right) + \dots$$

For example

$$g_{ij}^{(1)}(x^i) = -L^2 P_{ij}[g^{(0)}] = -\frac{L^2}{d-2} \left(R_{ij}[g^{(0)}] - \frac{g^{(0)}_{ij}}{2(d-1)} R[g^{(0)}] \right)$$

Induced gravity action

FG expansion into the action, and integrating out the z direction

$$I_{\text{eff}} = \frac{1}{16\pi G_{\text{eff}}} \int d^d x \sqrt{-\tilde{g}} \left[\frac{(d-1)(d-2)}{\ell_{\text{eff}}^2} + \tilde{R}(\tilde{g}) \right] \\ + \frac{1}{16\pi G_{\text{RS}}} \int d^d x \sqrt{-\tilde{g}} \left[\frac{L^2}{(d-4)(d-2)} \left(\tilde{R}^{ij} \tilde{R}_{ij} - \frac{d}{4(d-1)} \tilde{R}^2 \right) + \dots \right]$$

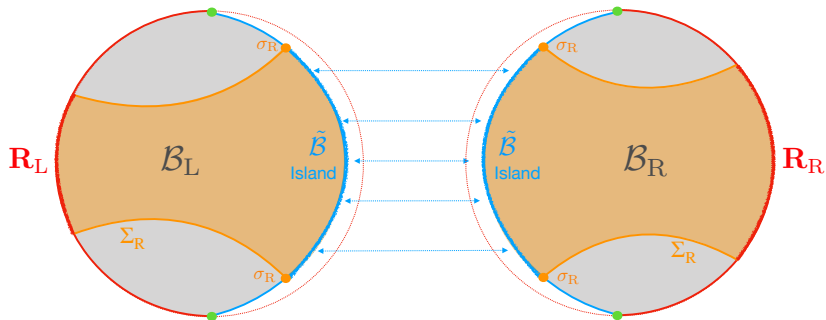
where

$$\frac{1}{G_{\text{eff}}} = \frac{1}{G_{\text{RS}}} = \frac{2L}{(d-2)G_{\text{bulk}}}, \quad \frac{L^2}{\ell_{\text{eff}}^2} \approx \frac{z_B^2}{L^2} \approx 2 - \frac{8\pi G_{\text{bulk}} L T_o}{d-1}.$$

Locality and small curvature limit

Induced action is a series in $L^2 \times \tilde{R} \sim L^2/\ell_{\text{eff}}^2 \sim z_B^2/L^2$

Brane close to asymptotic boundary $z_B/L \ll 1$

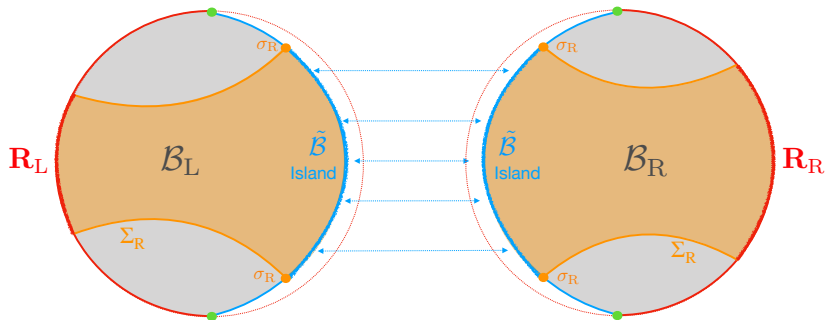


Natural expansion parameter z_B/L

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Brane close to asymptotic boundary $z_B/L \ll 1$



Natural expansion parameter z_B/L

Note: lengths on the brane scale as L/z_B

Intrinsic gravitational action

Add a gravitational action to the brane

$$I_{\text{brane}} = -(T_0 - \Delta T) \int d^d x \sqrt{-\tilde{g}} + \frac{1}{16\pi G_{\text{brane}}} \int d^d x \sqrt{-\tilde{g}} \tilde{R},$$

where ΔT ensures the location of the brane doesn't change.

FG expansion + z integration lead to same induced gravity action, but now

$$\frac{1}{G_{\text{eff}}} = \frac{1}{G_{\text{RS}}} + \frac{1}{G_{\text{brane}}} = \frac{2L}{(d-2)G_{\text{bulk}}} + \frac{1}{G_{\text{brane}}}$$

For $d=2$, I_{eff} is non-local $\sim \tilde{R} \log \tilde{R} \sim$ on-shell Polyakov action

For $d=2$, I_{EH} is topological, so one can instead add a I_{JT} on the brane

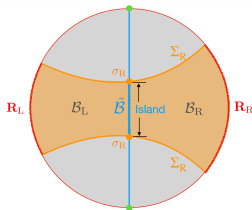
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Entanglement entropy of the island

Holographic entanglement entropy

$$S(R) = \text{ext}_{\partial\Sigma_R=\partial R} \frac{A(\Sigma_R)}{4G_{\text{bulk}}}$$



Applying a Fefferman-Graham expansion

$$x^i(z, \sigma^a) = x^i{}^{(0)}(\sigma^a) + \frac{z^2}{L^2} x^i{}^{(1)}(\sigma^a) + \dots, \quad x^i{}^{(1)} = \frac{L}{2(d-2)} K n^i{}^{(0)}$$

about the brane on holographic entanglement entropy

$$S(R) = \frac{A(\Sigma_R)}{4G_{\text{bulk}}} = UV(\partial R) + \frac{\tilde{A}(\sigma_R)}{4G_{\text{eff}}} + \dots$$

Entanglement entropy of the island

[Chen, Myers, Neuenfeld, Reyes, Sandor] and [JH, Myers, Ruan]

Including subleading terms, the contribution to entanglement entropy is the Wald-Dong entropy for the induced action on the brane

$$S(R) = UV(\partial R) + S_{WD}(\sigma_R) + \dots$$

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- Since RT surface Σ_R extremizes S , it follows that QES σ_R extremizes S_{WD}

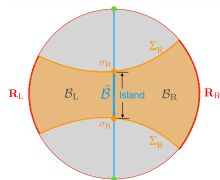
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$$S(R) = UV(\partial R) + S_{WD}(\sigma_R) + \dots$$

- Since RT surface Σ_R extremizes S , it follows that QES σ_R extremizes S_{WD}
- $UV(\partial R) \sim S_{WD}(\partial R) \sim UV$ entanglement across ∂R
- $S_{WD}(\sigma_R) \sim \mathcal{O}(z_B/L)$ entanglement across QES σ_R
- Recall z_B/L is locality scale of the brane

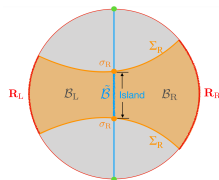


Complexity of the island

[JH, Myers, Ruan]

Subregion complexity=volume

$$C_V(R) = \max_{\partial\mathcal{B}=R \cup \Sigma_R} \frac{V(\mathcal{B})}{G_{\text{bulk}} \ell}$$

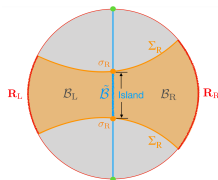


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Applying a Fefferman-Graham expansion about the brane on holographic complexity=volume

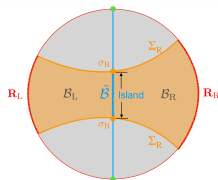
$$C_V(R) = \frac{V(\mathcal{B})}{G_{\text{bulk}}\ell} = UV(R) + \frac{\tilde{V}(\tilde{\mathcal{B}})}{G_{\text{eff}}\ell} + \dots$$

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$$C_V(R) = \frac{V(\mathcal{B})}{G_{\text{bulk}}\ell} = UV(R) + \frac{\tilde{V}(\tilde{\mathcal{B}})}{G_{\text{eff}}\ell} + \dots$$

UV divergences cancel when considering “mutual complexity”

$$C_V(R_L) + C_V(R_R) - C_V(R) = \frac{\tilde{V}(\tilde{\mathcal{B}})}{G_{\text{eff}}\ell} + \dots$$

Complexity of the island

[JH, Myers, Ruan]

Including subleading terms lead to a “generalized volume”

$$C_V(R) = UV(R) + \frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}} l} + \dots$$

Complexity of the island

[JH, Myers, Ruan]

Including subleading terms lead to a “generalized volume”

$$C_V(R) = UV(R) + \frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}} l} + \dots$$

- \widetilde{W} derived from the action in parallel to S_{WD}

Complexity of the island

[JH, Myers, Ruan]

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- Since \mathcal{B} maximizes C_V , it follows that the island $\tilde{\mathcal{B}}$ maximizes \widetilde{W}

Complexity of the island

[JH, Myers, Ruan]

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- $UV(R) = \frac{\widetilde{W}(R)}{G_{\text{eff}}\ell} \sim$ complexity of UV entanglement structure in R

Complexity of the island

[JH, Myers, Ruan]

Including subleading terms lead to a “generalized volume”

$$C_V(R) = UV(R) + \frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}} \ell} + \dots$$

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- Since \mathcal{B} maximizes C_V , it follows that the island $\tilde{\mathcal{B}}$ maximizes \widetilde{W}
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- $\frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}} \ell} \sim$ complexity of $\mathcal{O}(z_B/L)$ entanglement structure in island $\tilde{\mathcal{B}}$

Complexity of the island

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- Unlike EE, subregion complexity is discontinuous at the island transition

Complexity of the island

[JH, Myers, Ruan]

Including subleading terms lead to a “generalized volume”

$$C_V(R) = UV(R) + \frac{\widetilde{W}(\widetilde{\mathcal{B}})}{G_{\text{eff}}\ell} + \dots$$

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- $\frac{\widetilde{W}(\widetilde{\mathcal{B}})}{G_{\text{eff}}\ell} \sim$ complexity of $\mathcal{O}(z_B/L)$ entanglement structure in island $\widetilde{\mathcal{B}}$
- Unlike EE, subregion complexity is discontinuous at the island transition
- The complexity gained in the island phase is that of the most complex island anchored at the QES

Gravity on the brane

[JH, Myers, Ruan]

We can add an intrinsic gravitational term to the brane action

$$I_{\text{brane}} = -(T_o - \Delta T) \int d^d x \sqrt{-\tilde{g}} + \frac{1}{16\pi G_{\text{brane}}} \int d^d x \sqrt{-\tilde{g}} \tilde{R}$$

Because of the intrinsic gravity on the brane, C_V gets an extra brane contribution

$$C_V(R) = \max_{\partial\mathcal{B}=R \cup \Sigma_R} \frac{V(\mathcal{B})}{G_{\text{bulk}} \ell} + \frac{\tilde{V}(\tilde{\mathcal{B}})}{G_{\text{brane}} \ell'}$$

The same result

$$C_V(R) = UV(R) + \frac{\tilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}} \ell} + \dots$$

where

$$\frac{1}{G_{\text{eff}}} = \frac{2L}{(d-2)G_{\text{bulk}}} + \frac{1}{G_{\text{brane}}}$$

Same for entropy

Generalized CV

And higher curvature bulk gravity

Propose a generalized C_V

$$C_V(R) = \max_{\partial\mathcal{B}=R\cup\Sigma_R} \frac{W(\mathcal{B})}{G_{\text{bulk}}\ell} + \frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{brane}}\ell'}$$

and verify

$$C_V(R) = UV(R) + \frac{\widetilde{W}(\tilde{\mathcal{B}})}{G_{\text{eff}}\ell} + \dots$$

for Gauss-Bonnet and $F(R)$ gravity

$$\lambda_{GB} (R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma} - 4R_{\mu\nu}R^{\mu\nu} + R^2) , \quad F(R)$$

Closing comments

Entanglement entropy and complexity of island

- (Sub)leading contribution in the island phase: Wald entropy/complexity of island
- Scales with locality of induced gravity in the brane
- Double holography, what about more generally?

Discontinuity of subregion complexity

- General in holography
- Discontinuity of EW
- Large N effect

Outline

- 1 Entanglement entropy and complexity
- 2 Entanglement and complexity of islands
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 - Entanglement velocity and Lloyd bounds
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- 4 Conclusion

Bounds from entanglement and complexity

In braneworld models, the entanglement/complexity surfaces can sometimes intersect the branes

For branes with intrinsic gravity, there is an additional contact term that should be included in the entanglement entropy/complexity proposals

The gravitational couplings on the brane can influence the behaviour of entanglement/complexity

Bounds on behaviour for entanglement/complexity can limit the allowed values of gravitational couplings on the brane

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Entanglement velocity bound

[Afkhami-Jeddi, Hartman]

Normalized rate of entanglement growth

$$v_E \equiv \frac{\partial_t S(A)}{s_{th}(\beta)|\partial A|},$$

For 2d CFTs, states with uniform energy density, the entanglement velocity is upper bounded

$$|v_E| \leq \coth \frac{\pi|A|}{\beta}.$$

In the thermodynamic limit $|A| \gg \beta$,

$$|v_E| \leq 1.$$

Lloyd bound

“Energy limits speed of computation” [Lloyd]

Consider the “ultimate laptop”

- Average temperature E
- Bound on rate of logical operations
- Operations/second $\leq \frac{2E}{\pi\hbar}$

Bound on complexity growth

- $\frac{dC}{dt} \leq \frac{2E}{\pi\hbar}$
- Saturated for holographic complexity of black holes
- Black holes as “ultimate computers”

Outline

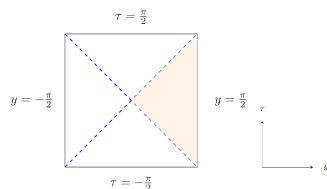
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Black hole + ETW brane

[Hartman, Maldacena], [Takayanagi] and [Fujita, Takayanagi, Tonni]

Single sided BH microstate geometry

- Double sided planar black hole
- Capped by ETW brane on one side
- BH microstate with semiclassical geometry



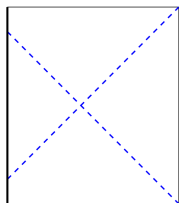
Total action

$$I = \frac{1}{16\pi G_N} \left[\int_{\text{bulk}} d^{d+1}x \sqrt{-g} \left(R + \frac{d(d-1)}{L^2} \right) + 2 \int_{\partial \text{AdS}} d^d x \sqrt{-h} K \right. \\ \left. + 2 \int_{\text{brane}} d^d x \sqrt{-h} \left(K - \frac{T_0}{L^{d-1}} \right) \right].$$

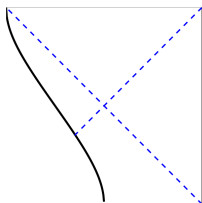
Black hole + ETW brane

[Lee, Shukla, Neuenfeld]

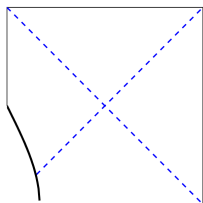
Israel junction conditions determine location of the brane



(a) Subcritical: $0 \leq T_0 < 1$



(b) Critical: $T_0 = 1$



(c) Supercritical: $T_0 > 1$

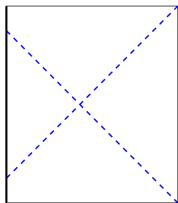
In addition, we can add gravity action on the brane

$$I_{\text{brane}} = \frac{1}{9\pi G_N} \int_{\text{brane}} d^d x \sqrt{-h} \left(K - \frac{T}{L^{d-1}} \right) + \frac{1}{16\pi G_N^{\text{brane}}} \int_{\text{brane}} d^d x \sqrt{-h} R^{\text{brane}}$$

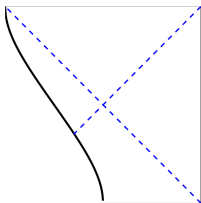
Black hole + ETW brane

[Lee, Shukla, Neuenfeld]

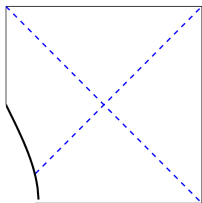
Israel junction conditions determine location of the brane



(a) Subcritical: $0 \leq T_0 < 1$



(b) Critical: $T_0 = 1$



(c) Supercritical: $T_0 > 1$

In addition, we can add intrinsic JT gravity on the brane

$$I_{\text{JT}} = \frac{1}{16\pi G_N^{\text{brane}}} \int d^2x \sqrt{-h} \varphi \left(R^{\text{brane}} - 2\Lambda^{\text{brane}} \right) + \frac{1}{16\pi G_N^{\text{brane}}} \int d^2x \sqrt{-h} \Phi_0 R^{\text{brane}} .$$

Brane location and dilaton profile

[Lee, Shukla, Neuenfeld]

The location of the brane is determined by the cosmological constant

$$-\frac{1}{L^2} \cos^2 y_{\text{brane}} = \Lambda^{\text{brane}},$$

and the dilaton profile is found by varying the metric on the brane

$$\varphi(\tau_{\text{brane}}) = \varphi_0 + \varphi_1 \sin \tau_{\text{brane}},$$

where $\varphi_0 = G_N^{\text{brane}} K / 2G_N \Lambda^{\text{brane}}$ and φ_1 are constants.

Brane location and dilaton profile

[Lee, Shukla, Neuenfeld]

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There are bounds on the JT coupling φ_1 due to entanglement and complexity growth bounds.

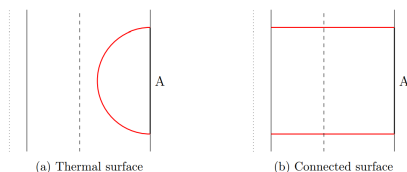
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Entanglement bounds

[Lee, Shukla, Neuenfeld]

Consider an interval A large enough that the connected HRT surface still dominates.



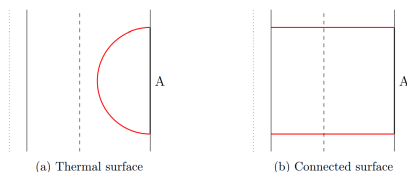
Include contact term due to intrinsic gravity on the brane

$$S(A) = \frac{A(\Sigma)}{4G_N} + \frac{\phi(\Sigma \cap \text{brane})}{4G_N^{\text{brane}}}.$$

Entanglement bounds

[Lee, Shukla, Neuenfeld]

Consider an interval A large enough that the connected HRT surface still dominates.



Include contact term due to intrinsic gravity on the brane

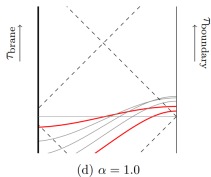
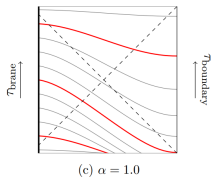
$$S(A) = \frac{A(\Sigma)}{4G_N} + \frac{\phi(\Sigma \cap \text{brane})}{4G_N^{\text{brane}}}.$$

Impose the upper bound on entanglement velocity $|v_E| \leq 1$.

Entanglement bounds

Two bounds on JT coupling $\alpha \equiv \frac{G_N \varphi_1}{G_N^{\text{brane}} L}$

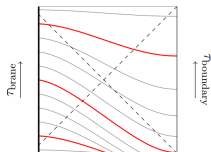
- Rapid growth of entanglement
- Discontinuous jump in entanglement



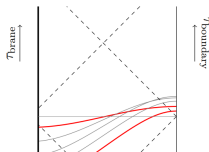
Entanglement bounds

Two bounds on JT coupling $\alpha \equiv \frac{G_N \varphi_1}{G_N^{\text{brane}} L}$

- Rapid growth of entanglement
- Discontinuous jump in entanglement



(c) $\alpha = 1.0$



(d) $\alpha = 1.0$

First bound:

$$|\alpha| \leq \frac{1}{1 + \sin y_{\text{brane}}}.$$

Second bound:

$$|\alpha| \leq \frac{1}{1 - \sin y_{\text{brane}}}.$$

Entanglement bounds

Two bounds on JT coupling $\alpha \equiv \frac{G_N \varphi_1}{G_N^{\text{brane}} L}$

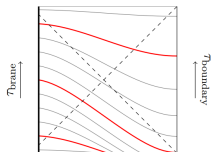
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First bound:

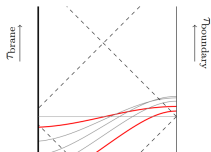
$$|\alpha| \leq \frac{1}{1 + \sin y_{\text{brane}}}.$$

Second bound:

$$|\alpha| \leq \frac{1}{1 - \sin y_{\text{brane}}}.$$



(c) $\alpha = 1.0$



(d) $\alpha = 1.0$

Entanglement velocity asymptotes to its upper bound

$$v_E = 1 - 2\mathcal{K}_1 e^{\frac{-4\pi t}{\beta}} + \dots,$$

where $\mathcal{K}_1 = \frac{1 - \alpha - \alpha \sin y_{\text{brane}}}{1 + \alpha - \alpha \sin y_{\text{brane}}}$.

Complexity bounds

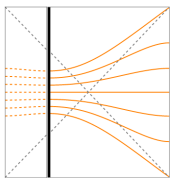
[Aguilar-Gutierrez, Craps, JH, Shulka, Khramtzov, Knysh] (w.i.p)

Recall the Lloyd bound on complexity

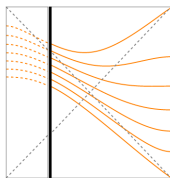
$$\frac{dC}{dt} \leq \frac{2M}{\pi\hbar}$$

Compute the growth rate of complexity=volume. Include contact term due to intrinsic gravity on the brane

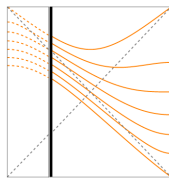
$$C_V(R) = \frac{V(\mathcal{B})}{G_N L} + \frac{1}{G_N^{\text{brane}} L} \int_{\mathcal{B} \cap \text{brane}} \phi.$$



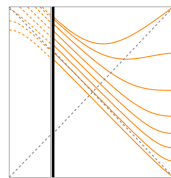
(a) $\alpha = 0$



(b) $\alpha = 1$



(c) $\alpha = 1.42$



(d) $\alpha = 5$

Complexity bounds

Tuning the Lloyd bound

- Definitions of holographic complexity are ambiguous
- Universal late time growth rate
- Proportional to M
- Intuition from entanglement velocity, and BHs as “ultimate computers”

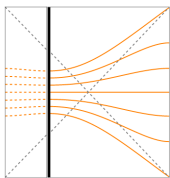
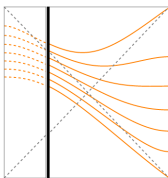
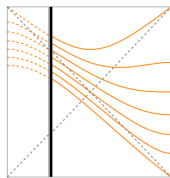
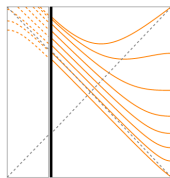
Use this universal late time value as the Lloyd bound.

Complexity bounds

Tuning the Lloyd bound

- Definitions of holographic complexity are ambiguous
- Universal late time growth rate
- Proportional to M
- Intuition from entanglement velocity, and BHs as “ultimate computers”

Use this universal late time value as the Lloyd bound. With this tuning, complexity surface projected to fall into singularity = violation of Lloyd bound

(a) $\alpha = 0$ (b) $\alpha = 1$ (c) $\alpha = 1.42$ (d) $\alpha = 5$

Complexity bounds

The bound on the JT coupling $\alpha \equiv \frac{G_N \varphi_1}{G_N^{\text{brane}} L}$ is

$$|\alpha| \leq -\frac{\tan y_{\text{brane}}}{\cos y_{\text{brane}}},$$

where, recall

$$-\frac{1}{L^2} \cos^2 y_{\text{brane}} = \Lambda^{\text{brane}}.$$

Entanglement and complexity bounds

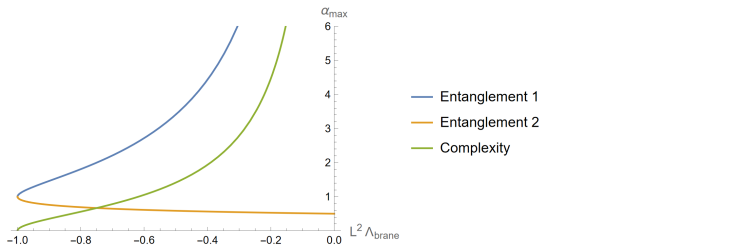
The bounds on the JT coupling α from entanglement velocity are

$$|\alpha| \leq \frac{1}{1 - \sin y_{\text{brane}}}, \quad |\alpha| \leq \frac{1}{1 + \sin y_{\text{brane}}}.$$

The bounds on α from Lloyd bound is

$$|\alpha| \leq -\frac{\tan y_{\text{brane}}}{\cos y_{\text{brane}}},$$

where $-\frac{1}{L^2} \cos^2 y_{\text{brane}} = \Lambda^{\text{brane}}$.



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Closing comments

Entanglement entropy and complexity of island

- Double holography, what about more generally?

Discontinuity of subregion complexity

- General in holography
- Discontinuity of EW

Entanglement and complexity bounds

- Currently no strict bound on entanglement velocity in higher dimensions
- There are late time bounds
- Could apply to other braneworld models

Open questions

Entanglement and complexity of islands

- C_A , CV2.0 of the island
- C =anything of the island
- Higher order corrections to S_{WD} , W
- W from first law of causal diamonds (Bueno et al.)

Entanglement and complexity bounds

- Bounds from complexity in higher dimension
- Bounds on entanglement velocity in higher dimension
- Bounds from other holographic complexities
- Entanglement and complexity bounds in other braneworld models