

# Rapidity gap dynamics: color fluctuations vs knockout mechanism

**Mark Strikman, PSU**

**In collaboration with L.Frankfurt, V.Guzey, A. Stasto, W.Lee, M.Zhalov**

Reports on Progress in Physics REVIEW

Selected topics in diffraction with protons and nuclei: past, present, and future

L Frankfurt<sup>1,2</sup>, V Guzey<sup>4,3</sup>, A Stasto<sup>2</sup> and M Strikman<sup>2</sup>Published 26 October 2022 • © 2022 IOP Publishing Ltd

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To describe soft diffraction Good and Walker assumed that the state of an energetic

incident hadron  $|\Psi\rangle$   $|\Psi\rangle = \sum_k c_k |\Psi_k\rangle$  ,

$$\text{Im}T|\Psi_k\rangle = t_k|\Psi_k\rangle ,$$

$$\sum_k |c_k|^2 = 1 .$$

$|\Psi_k\rangle$  interact with the target with different cross sections  $\sigma$  and orthogonal

$$\left(\frac{d\sigma}{dt}\right)_{t=0}^{\text{diff}} = \frac{1}{16\pi} \sum_k |\langle \Psi_k | \text{Im}T | \Psi \rangle|^2 = \frac{1}{16\pi} \sum_k |c_k|^2 t_k^2 = \frac{1}{16\pi} \langle \sigma^2 \rangle$$

$$\left(\frac{d\sigma}{dt}\right)_{t=0}^{\text{el}} = \frac{1}{16\pi} |\langle \Psi | \text{Im}T | \Psi \rangle|^2 = \frac{1}{16\pi} \left( \sum_k |c_k|^2 t_k \right)^2 = \frac{1}{16\pi} \langle \sigma \rangle^2 .$$

**Inelastic diffractions is related to fluctuations of cross section**

**Problem GOOD & WALKER** LOGIC TO BE APPLICABLE FOR A RANGE OF T , IT ISN  
NECESARY TO HAVE THE SAME EIGENSTATES FOR DIFFERENT T. I

**Doubtful** In pQCD DIPOLES OF DIFFERENT SIZE ARE EIGENSTATES FOR T=0, BUT NOT FOR FINITE T.

**Nonperturbative example — deuteron +p → d+p & pn +p.**

Impulse approximation  
No fluctuations

the condition of orthogonality of the wave functions  
of the continuum and bound states leads to

$$\left. \frac{d\sigma(D + h \rightarrow "pn" + h)}{dt} \right|_{t=0} = 0,$$

Consistent with GW

**But contrary to GW expectations at finite t**

$$\left. \frac{d\sigma(D + h \rightarrow "pn" + h)}{dt} \right|_{\text{incoh}} = 2 \frac{d\sigma(N + h \rightarrow N + h)}{dt} (1 - F_D^2(4t))$$

$$\frac{d\sigma(D + h \rightarrow D + h)}{dt} = 4 \frac{d\sigma(N + h \rightarrow N + h)}{dt} F_D^2(t),$$

**Recently renewed interest to Inelastic diffraction in  $\gamma + p (A) \rightarrow J/\psi$  (leading dijet) + gap + Y**

Rapidity gap In particular it was suggested by Heikki Mäntysaari, Björn Schenke that the data at  $-t < 2 \text{ GeV}^2$  are dominated by color fluctuations. Hot spots, etc.

Prompted us to look again on these processes - interplay of different mechanisms which relative role depends on  $t$ .

# Inelastic diffraction in $\gamma + p (A) \rightarrow J/\psi$ (leading dijet) + gap + Y

## Three regimes

- **t=0** - color fluctuations in nucleons variance of

$$\sigma_{diff} / \sigma_{el} = \text{variance of gluon density at given } x \text{ (color fluctuations)}$$

Frankfurt et al

- **-t > 0.3 ÷ 0.5 GeV<sup>2</sup>** elastic scattering of a small dipole off gluons & quarks

for smaller t this mechanism is suppressed by factor

$$R = 1 - \left( \frac{1}{1 - t/M^2} \right)^4$$

**M<sup>2</sup> = 1 GeV<sup>2</sup>**

- **0.1 < -t < 0.3 ÷ 0.5 GeV<sup>2</sup>** interplay of these two mechanisms  
plus spinflip

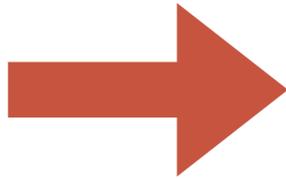
- $t=0$  Processes where  $\omega_g$  is coupled to two gluon ladder

Logic is a combination of QCD factorization theorem for exclusive processes and Good Walker's

$$G(x, Q^2) = \sum_n |a_n|^2 G(x, Q^2|n) \equiv \langle G \rangle$$

$$\left. \frac{d\sigma_{elastic}}{dt} \right|_{t=0} \propto \left[ \sum_n |a_n|^2 G(x, Q^2|n) \right]^2 \equiv \langle G \rangle^2$$

$$\left. \frac{d\sigma_{diffractive}}{dt} \right|_{t=0} \propto \sum_n |a_n|^2 [G(x, Q^2|n)]^2 \equiv \langle G^2 \rangle$$

  $\omega_g \equiv \frac{\langle G^2 \rangle - \langle G \rangle^2}{\langle G \rangle^2} = \left[ \frac{d\sigma_{inelastic}}{dt} / \frac{d\sigma_{elastic}}{dt} \right]_{t=0}$  Frankfurt et al

$$\sim 0.15 - 0.2$$

**Perturbative Pomeron: what is energy dependence cross section in vacuum channel ?**

Problem for the study - two large parameters  $\ln Q^2$ , and  $\ln 1/x$ .

DIS - both parameters enter (DGLAP); BGKL - only  $\ln 1/x$  (scattering of two small dipoles)

BFKL elastic amplitude  $f(s) = (s/s_0)^{1 + \omega_P}$

$$\omega_{/P} = a_1 \alpha_S - a_2 \alpha_S^2 + \dots$$

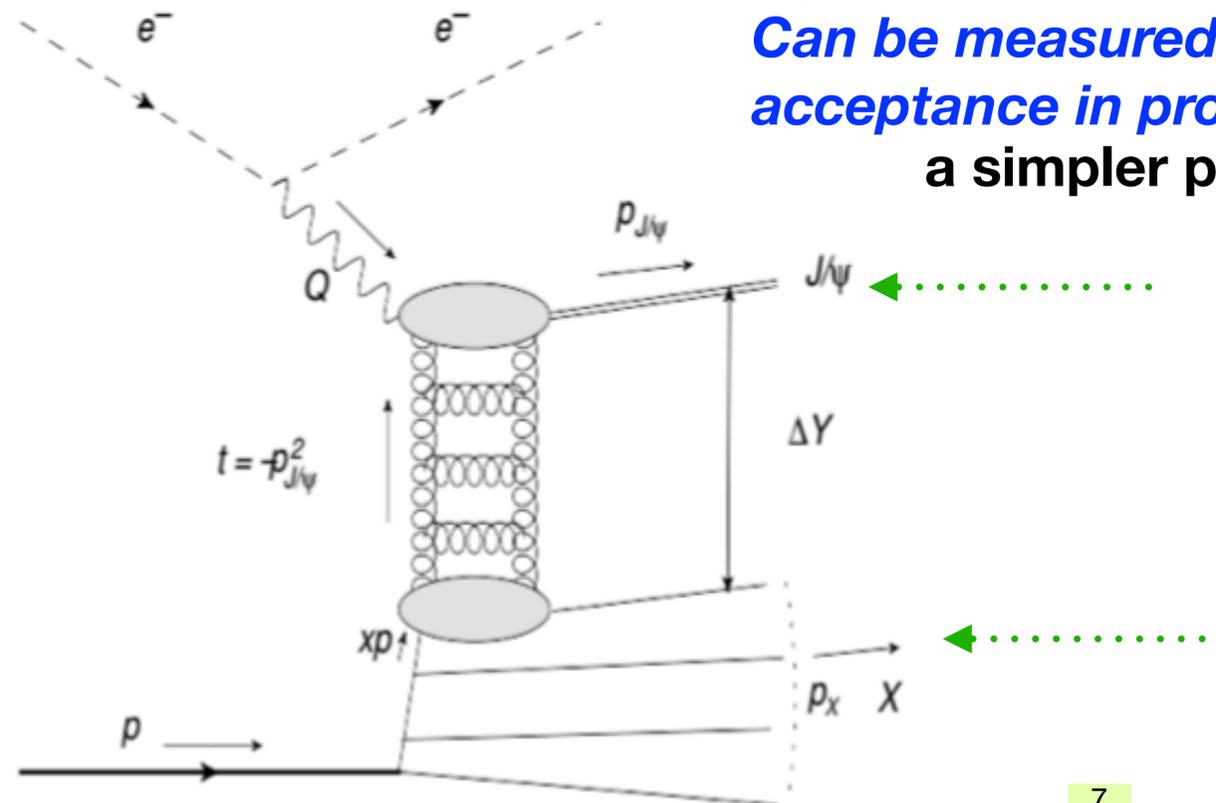
leading log  $\omega_{/P} \sim 0.5 \div 0.8$ , NLO  $\sim 0.1$ , resummation  $\sim 0.25$

Main reason for small values of  $\omega_{/P}$  energy conservation

*Promising direction: Rapidity gaps at large  $t$  for  $J/\psi$  production - squeezing from both ends.*

*Can be measured in UPC (pA) if good acceptance in proton region*

a simpler process than Mueller and Tung dijet

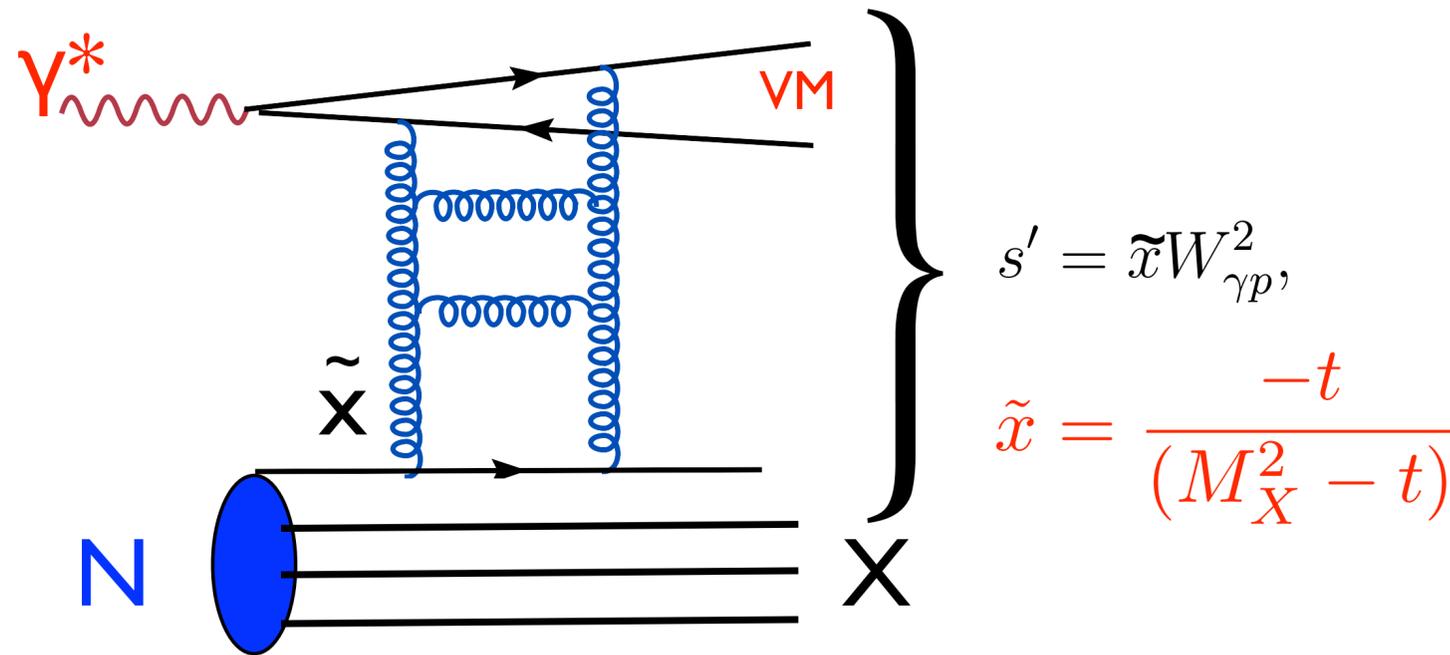


**rapidity gap**

*Parton knockout. mechanism*

**fixed x**

elementary reaction scattering of projectile off a parton of the target at large  $t$  belongs to a class of reactions with hard white exchange in t-channel



FS 89, FS95,  
Mueller & Tung 91  
Forshaw & Ryskin 95

best way to measure of the strength of inelastic interactions of small dipole in the processes initiated by elastic small dipole - parton scattering. In HI via UPC feasible for at  $[s']^{1/2} = 20 \text{ GeV} - 100 \text{ GeV}$  at the LHC

The choice of large  $t$  ensures several important simplifications:

\* the parton ladder mediating quasielastic scattering is attached to the projectile via two gluons.

\*\* attachment of the ladder to two partons of the target is strongly suppressed.

\*\*\* small transverse size  $d_{q\bar{q}} \propto 1/\sqrt{-t} \sim 0.15\text{fm}$  for  $J/\psi$  for  $-t \sim m_{J/\psi}^2$

$$\frac{d\sigma_{\gamma+p \rightarrow V+X}}{dt d\tilde{x}} = \frac{d\sigma_{\gamma+quark \rightarrow V+quark}}{dt} \left[ \frac{81}{16} g_p^{31}(\tilde{x}, t) + \sum_i (q_p^i(\tilde{x}, t) + \bar{q}_p^i(\tilde{x}, t)) \right]$$

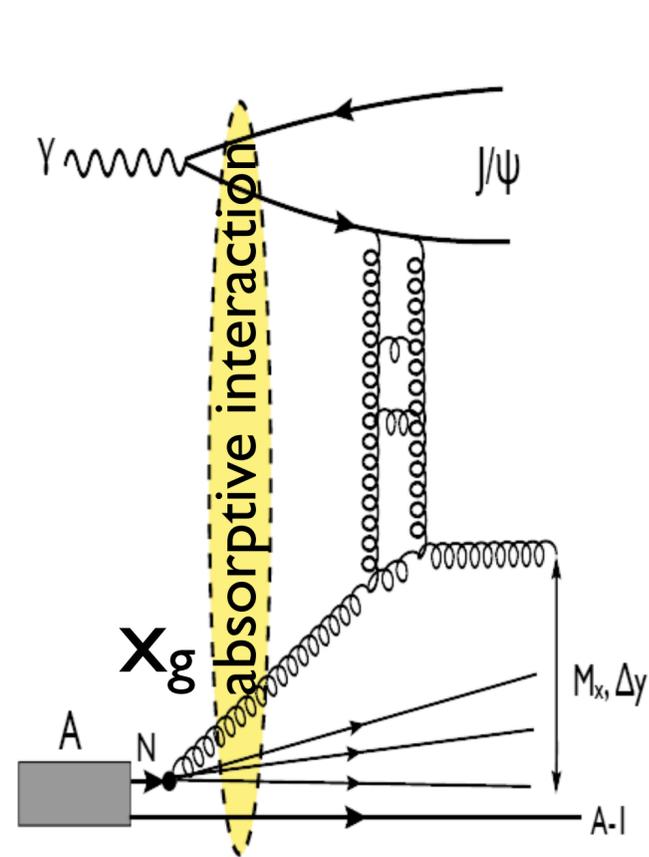
$\exp(2\omega_p \Delta Y)$

resummation predicts a huge effect - between  $\Delta Y = 2$  and  $\Delta Y = 4$   $\sigma$  is expected to increase by a factor of 3 !!!

if EIC would have a detector with high acceptance in the nucleon fragmentation region. At LHC much larger  $\Delta Y$  can be reached  $\rightarrow$  even larger effect



Complementary to coherent  $J/\psi$ .

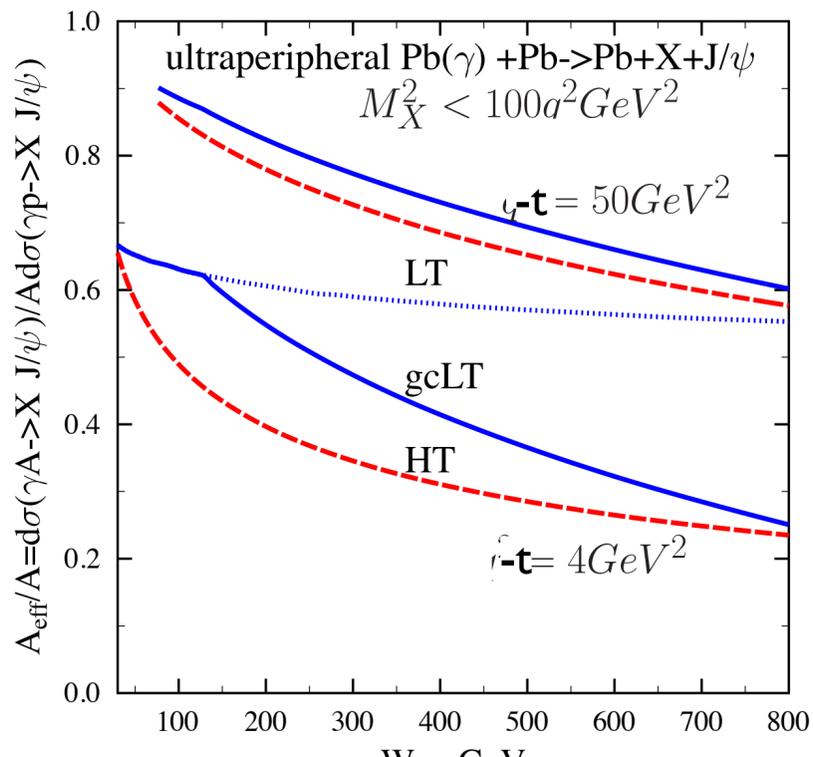


$s' = \tilde{x} W_{\gamma N}^2$

- ✱ Tracks dipole though  $\sim 10$  fm of nuclear matter
- ✱ Allows to measure dipole size as a function of  $q^2$
- ✱ Can reach maximal  $W$  of LHeC

Estimate of dipole size for  $q^2=0$

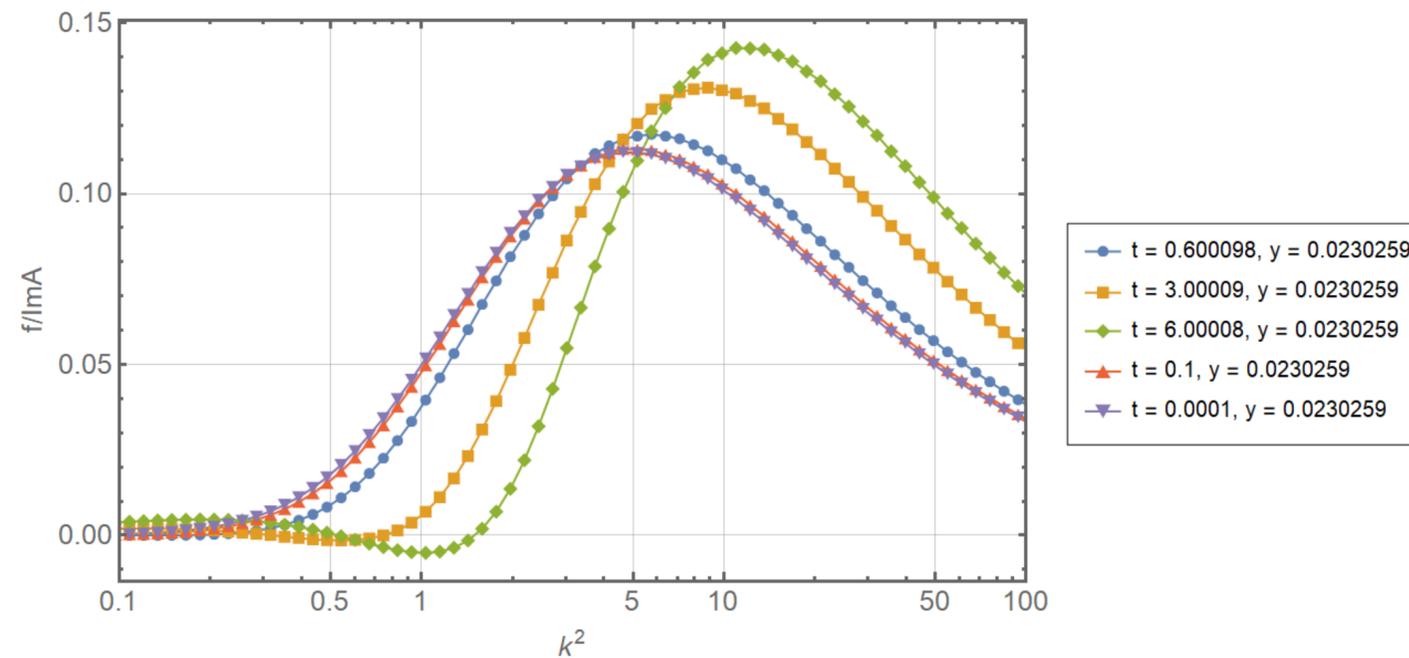
$d^2(-t)/d^2(0) \approx (1 - t/4m_c^2)^{-1}$   
 ( $d_0 = .25\text{fm}, m_c = 1.5\text{GeV}$ )



Significant absorption is expected in the leading twist and higher twist models of dipole interaction. One can select  $x_g$  both in the LT shadowing region and above 0.01

Large t data described well by pQCD. Iwe it crazy to go ti t~0?

We calculated distribution in the two gluon ladder over transverse momenta of gluons



Conclusion: momenta typically above 1 GeV/c: k

Two contributions - two gluons are attached to the same parton or to two different partons.

No suppression

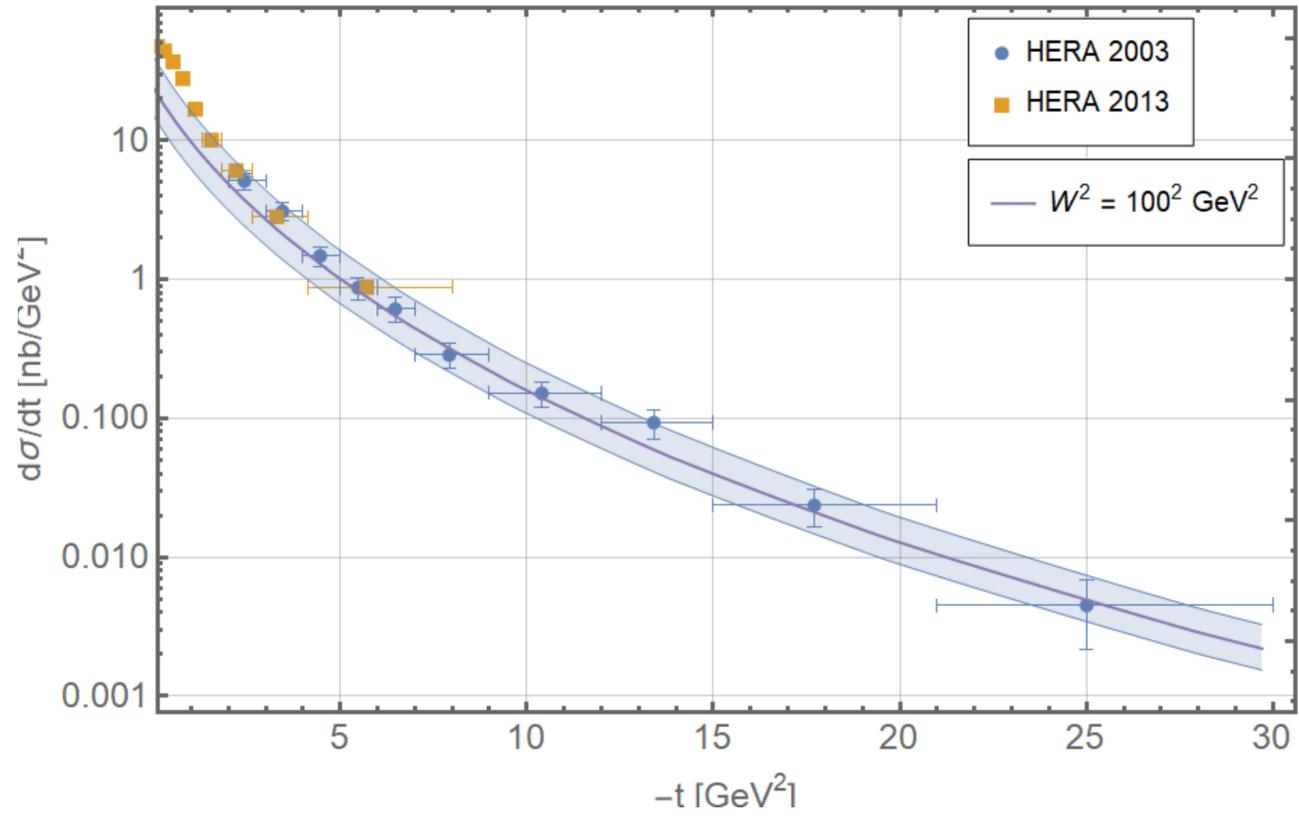
Large suppression.  $F_N^4(k^2)$

For  $t \rightarrow 0$  attachment to the same partons leads to cross section which is strongly suppressed - only elastic term survives (like in the deuteron example, logic similar to GW.)

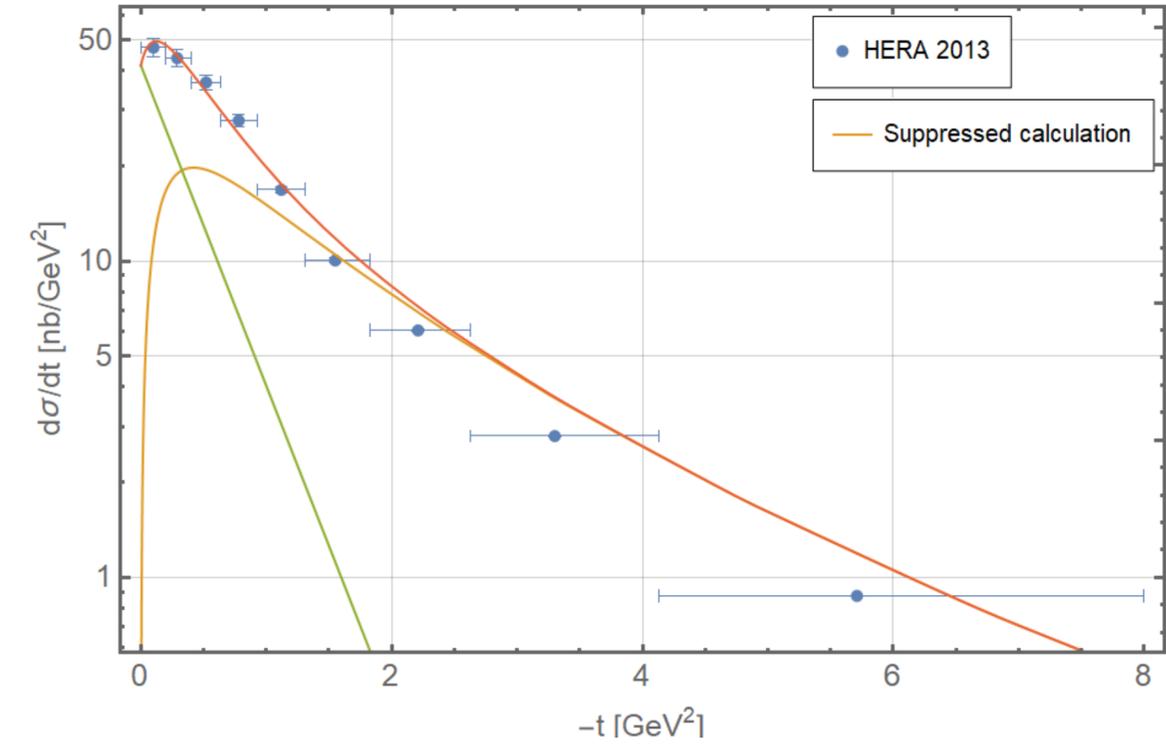
Suppression factor for each vertex - gluon form factor

$$R = 1 - \left( \frac{1}{1 - t/M^2} \right)^4$$

$$M \sim 1 \text{ GeV.}$$



Large t calculation extrapolated to small t



Suppression included.

Overall no need for gluon hot spots.  
Fluctuation cross section t-slope is comparable soft physics, in particular spin flip contribution

## CONCLUSIONS

*Gap physics with  $J/\psi$  promising direction for studying in ultra peripheral collisions*

NLO BFKL dynamics in a wide range of virtualities predicts fast energy dependence

$t=0$  inel/el tests onset of black regime

Good Walker model is justified (as a useful model for  $t$  close to 0)

Interesting nuclear effects

Easier to measure with improved rapidity  
acceptance of the LHC detectors





